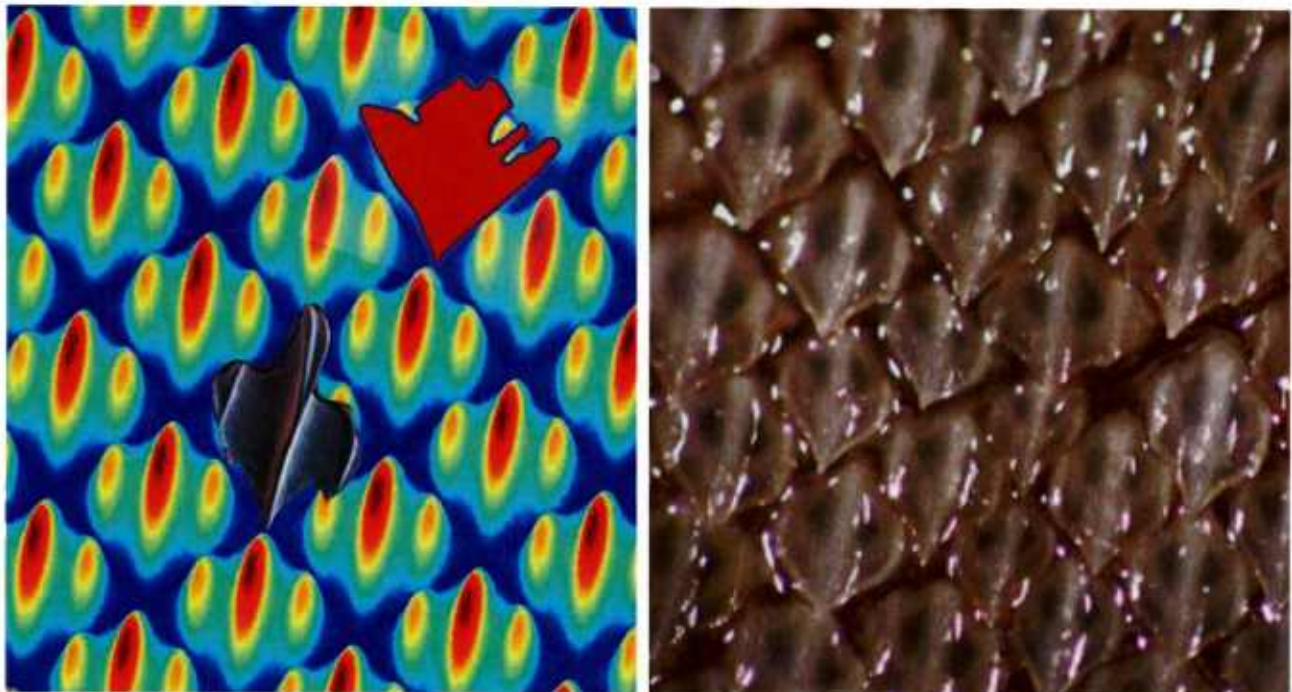


Llinas' Phase Reset Mechanism Delays the Onset of Chaos in Shark and Dolphin Wall Turbulence

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PREFACE

This report was prepared under NWA 100000802958, principal investigator Promode R. Bandyopadhyay (Code 851). The sponsoring activity is the Office of Naval Research (ONR-341, Thomas McKenna).

The technical reviewer for this report was James L. Dick (Code 8514).

The primary author gratefully acknowledges productive discussions with T. McKenna, R. Llinas, V. Makarenko, J. Simmons, and S. Roy. Also, the authors thank J. C. S. Meng, W. R. C. Phillips, and A M. Savill for their comments on the manuscript. AMH was sponsored by the ONR-33 Post-Doctoral Program.

Reviewed and Approved: 10 February 2014

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On the cover: *Left:* Model of the periodic distribution of the lateral diffusion of vorticity—compare this with the shark skin photograph on the right; upper red overlay is a model of the sub-harmonic map of the primary vorticity; lower overlay is a photograph of one dermal denticle of the Atlantic sharpnose shark (reference 36). Color codes: The fluctuating percentage of the real part of the lateral diffusion coefficient varies from 0.014 (red) to -0.01 (blue); subharmonic red is the maximum magnitude of the main component of the vorticity (0.6). *Right:* Photograph of skin of tiger shark (reference 37); compare this with the model on the left.

REPORT DOCUMENTATION PAGE				Form Approved OMB No. 0704-0188	
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1. REPORT DATE (DD-MM-YYYY) 10-02-2014		2. REPORT TYPE		3. DATES COVERED (From - To)	
4. TITLE AND SUBTITLE Llinas' Phase Reset Mechanism Delays the Onset of Chaos in Shark and Dolphin Wall Turbulence				5a. CONTRACT NUMBER	
				5b. GRANT NUMBER	
				5c. PROGRAM ELEMENT NUMBER	
6. AUTHOR(S) Promode R. Bandyopadhyay and Aren M. Hellum				5d. PROJECT NUMBER	
				5e. TASK NUMBER	
				5f. WORK UNIT NUMBER	
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) Naval Undersea Warfare Center Division 1176 Howell Street Newport, RI 02841-1708				8. PERFORMING ORGANIZATION REPORT NUMBER TR 12,154	
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES) Office of Naval Research One Liberty Center 875 North Randolph Street, Suite 1425 Arlington, VA 22203-2966				10. SPONSORING/MONITOR'S ACRONYM ONR	
				11. SPONSORING/MONITORING REPORT NUMBER	
12. DISTRIBUTION/AVAILABILITY STATEMENT Approved for public release; distribution is unlimited.					
13. SUPPLEMENTARY NOTES					
14. ABSTRACT Turbulent boundary layers (TBLs) develop in the flow over planetary surfaces, over the bodies of swimming animals (like sharks and dolphins), and over the hulls of ships and submarines. Evidence points to organization in turbulence production, despite the apparent randomness and lack of a unified theory. Here, the dynamical systems of neuroscience are used to develop a self-regulation theory to show for the first time how the skins of sharks and dolphins control the chaos of turbulence. To remove the apparent randomness, the turbulence production cycles are split into pre- and post-breakdown regions, both during the organized laminar-to-turbulent transition and in the TBL, which is chaotic. Stuart-Landau oscillators describe the instabilities of the surface-normal diffusion of vorticity, which is coupled to the instabilities' spanwise diffusion. Similar to Llinas' olivo-cerebellar temporal phase reset mechanism, the shark and dolphin skins impose combinations of vorticity diffusion templates internal to the oscillators and external microvibration or flow oscillations, synchronizing the oscillators laterally and thereby eliminating chaos. The theory—proving biological adaptation to environment—has potential application to tornado path management. Shark and dolphin skins are major targets of reverse engineering mechanisms of drag and noise reduction.					
15. SUBJECT TERMS <div style="display: flex; justify-content: space-between;"> Hydrodynamics Turbulent Boundary Layers Drag Reduction Transition Delay Swimming Animals </div> <div style="display: flex; justify-content: space-between;"> Biological Adaptation Sharks Dolphins Olivo-cerebellar Dynamics Chaos Control Dynamical Systems </div>					
16. SECURITY CLASSIFICATION OF:			17. LIMITATION OF ABSTRACT SAR	18. NUMBER OF PAGES 54	19a. NAME OF RESPONSIBLE PERSON Promode R. Bandyopadhyay
a. REPORT (U)	b. ABSTRACT (U)	c. THIS PAGE (U)			19b. TELEPHONE NUMBER (Include area code) 401-832-2712 email: promode.bandyopadhyay@navy.mil

20151007653

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LLINAS' PHASE RESET MECHANISM DELAYS THE ONSET OF CHAOS IN SHARK AND DOLPHIN WALL TURBULENCE

1. INTRODUCTION

Just as individual olivo-cerebellar neurons in animals appear noisy but *oscillatory* and their assemblage works reliably, in wall turbulence individual local states also appear noisy but *oscillatory* and laboratory experiments have demonstrated their global spatio-temporal organization. The skins of sharks and dolphins, which constantly contend with the seeming chaos of turbulence, also have organized patterns seemingly to control the chaos of wall turbulence. No analytical theory has been developed showing how this is done. Here, an analogy is drawn with olivo-cerebellar dynamics and the mechanisms of phase reset,¹ successfully splitting the pre- and post-breakdown parts of near-wall instabilities to develop a new spatio-temporal theory of wall turbulence and how sharks and dolphins use their skins to tame the chaos of turbulence.

The notable spatio-temporal observations of organized flow structures in wall turbulence are as follows. In transitional boundary layers, experimental discoveries of arrowhead turbulence spots,² arrayed and staggered shedding of hairpins,³ interpretations of hairpins⁴ and arrayed folds,⁵ and analytical theories^{6,7} showing organization have been reported. In turbulent boundary layers (TBLs), the following experimentally observed/deduced organized flow structures have been reported: inclined near-wall vortex pairs,⁸ streaks,⁹ bursting,¹⁰ hairpins,¹¹⁻¹⁴ pockets,¹⁵ large spatio-temporal structures with a characteristic upstream interface,^{12,16,17} and seemingly coupled large and small scales.¹⁸⁻²⁰

No single analytical theory of flow organization models simultaneously the features of transition, TBL, and control methods used by sharks and dolphins. It is not known how to separate randomness from organization. Also, it is not known what instability process (which is the key to control) takes place in the viscous region closest to the wall, even though this is an extensively observed and modeled region of the TBL. Separation of pre- and post-breakdown regions would reduce empiricism and help supplant post-dictive models of near-wall turbulence with predictive models.

This report offers a spatio-temporal theory of near-wall instabilities that give rise to organized flow structures, and it describes how to control the instabilities that produce these structures. It asserts that the apparent chaos of a TBL is the result of a group of self-regulating processes, similar to those apparent in biological processes.^{1,21-25} Significance is attached to two prior observations. First, the observation of a near-wall cyclic process²⁶ shows that the sublayer thickness initially rises while its growth rate slows, then suddenly undergoes breakdown into turbulence and the sublayer thickness falls. The present authors' interpretation is that *the sublayer undergoes a self-regulating growth-decay cycle; the vorticity accumulating in the sublayer during growth is abruptly liquidated into Strouhal arrays of hairpin vortices.*¹⁶ *The instant when this happens marks a point of bifurcation; the entire sublayer is always near*

bifurcation. The second prior observation is that high- and low-speed streaks converge and diverge.²⁷ The present authors' interpretation is that *the locations where the streaks converge and diverge are where the diffused sublayer vorticity is dislocating due to lateral diffusive coupling à la phase change in crystals*; hence, forms of the Ginzburg-Landau²⁸ (GL) theory of superconducting dislocation should phenomenologically apply.

The near-wall pre- and post-breakdown flows are split heuristically, choosing the complex SL oscillator description to describe the pre-breakdown instability process (see Methods in appendix A). Olivo-cerebellar control mechanisms¹ place the methods of chaos control on a common foundation. Assuming that a cellular flow has developed from a homogeneity due to end-effects, the theory shows how vorticity develops from a periodic to a chaotic state, and then—by dint of new principles of chaos control—how the periodic state is conserved, delaying the onset of chaos. Several disparate experimental and theoretical patterns are strikingly reproduced.

2. PRE- AND POST-BREAKDOWN REGIONS IN NEAR-WALL TURBULENCE ARE AUTONOMOUS

Figure 1a is a schematic delineating the initial and boundary condition-dependent pre- and post-bifurcation regimes of three-dimensional boundary layer growth and breakdown cycles. In the pre-bifurcation region, the vorticity is diffused but three-dimensional, whereas the bifurcation is characterized by an abrupt breakdown (it can be modeled as on/off switching); the diffused unsteady vorticity is liquidated down to 1.6 wall units²⁹ in Strouhal shedding of hairpin vortices.^{12,16} The pre-bifurcation self-regulating oscillations are modeled using complex SL oscillator equations. The turbulence states are described by oscillators of vorticities ω_z^+ , ω_x^+ , and viscous wall-layer thickness δ_s^+ , where the vorticities are coupled orthogonally to each other and diffusively in the spanwise direction. TBL measurements are used to calibrate the oscillators (Methods). No part of any state is filtered out as noise. Other oscillators such as ω_y^+ and pressure p^+ may also be considered. All results given are spatio-temporal. The present theory is a departure from the Direct Numerical Simulation (DNS) or Large Eddy Simulation (LES) solutions of the Navier-Stokes (NS) equations, or any turbulence modeling.

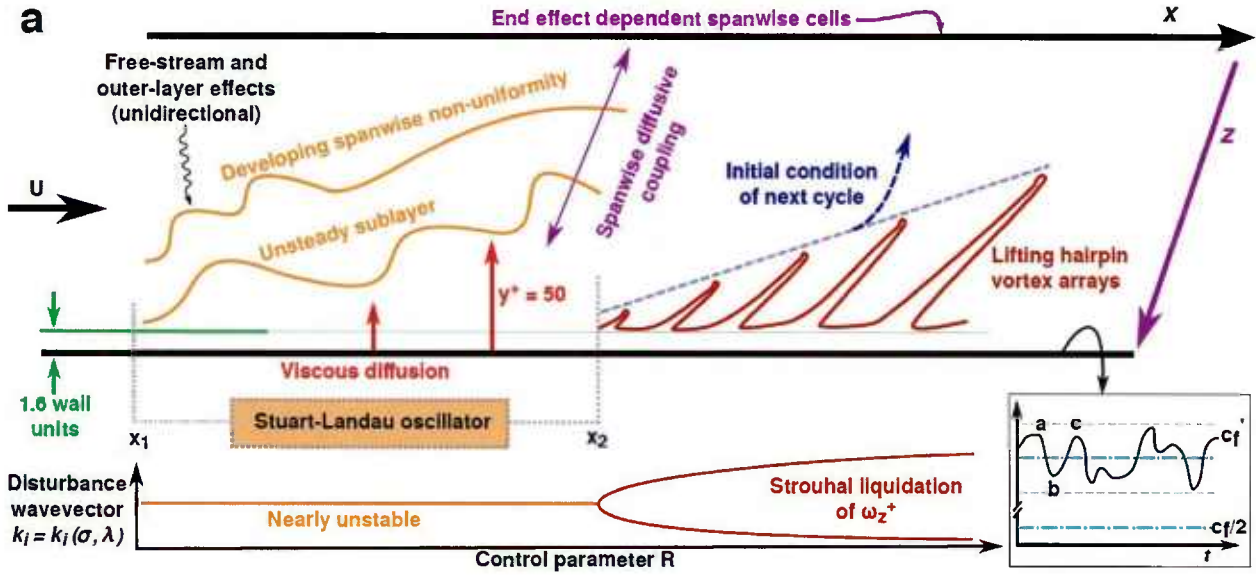


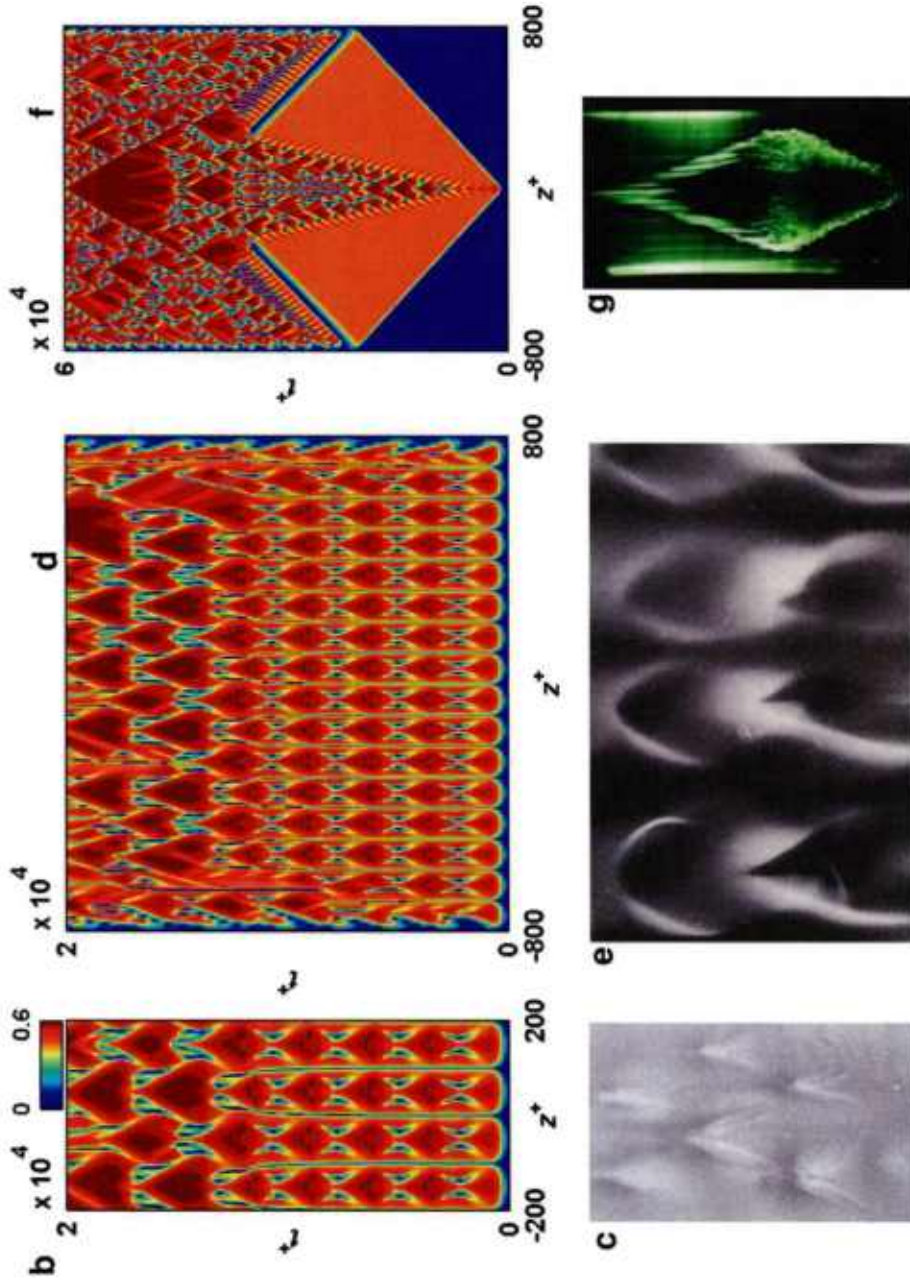
Figure 1a. Separation of Near-Wall TBL into Pre- and Post-Breakdown Regions (Regions of Molecular and Eddy Viscosity—Regions without and with Hairpin Vortices, Respectively)
(Inset: The turbulent part of skin friction (C_f') is self-regulated between limits (broken lines); the lower limit may drop to a minimum³⁰ of $C_f/2$.)

In the figure 1a inset, **a-b-c...** represents the decay and growth oscillation cycles of c_f' that oscillate, bounded between the dashed lines. If the process were linear, such as $\ddot{c}_f' + 2\zeta\omega\dot{c}_f' + \omega^2 c_f' = 0$, where ω is the frequency and ζ is the damping, the solution converges for $\zeta > 0$, diverges for $\zeta < 0$, and oscillates for $\zeta = 0$. For $\zeta = 0$, since the effect of the initial condition (IC) on the system c_f' is proportionate (no attenuation), the oscillations are large or small depending on the IC; furthermore, $\zeta < 0$. Instead, since the system output continues to

oscillate within bounds, the effect of the IC is being attenuated; that is, the process is nonlinear and self-correcting: $\ddot{c}_f' + 2f(c_f')\dot{c}_f' + \omega^2 c_f' = 0$, where the damping $f(c_f') = a_0 c_f'^2 - 2\zeta_0 \omega$ is negative when c_f' is small and positive when it is large. Here, a_0 is a constant, $(a_0, \zeta_0) > 0$, and a_0 , ζ_0 , and ω determine the shape of the limit cycle. Orr-Sommerfeld and GL or SL equations are examples of linear and nonlinear systems, respectively.

The post-bifurcation region is responsible for turbulence mixing characterized by eddy viscosity. The split of the TBL into pre- and post-bifurcation regimes reduces information overflow, allowing focus on the instability process. The present work focuses on the spatio-temporal instabilities of the pre-bifurcation regime (figures 1 and B-1—the latter in the Extended Data presented in appendix B), where the matters most important to chaos control and viscous drag reduction occur.

The arrayed and staggered distributions of vorticity during transition,³ known as K- and H-types,⁷ respectively, have been accurately reproduced. Time t^+ in theory and distance x^+ in smoke or dye flow visualization (which is a time history and is like a strip-chart record) are positive in opposite directions.³¹ Emmons' turbulence spot² is reproduced (figures 1b – 1g).



Figures 1b-1g. Comparison of Present Theories (z^+, x^+ ; c, e, g) in Respective Order $b: |\omega_z^+|$. Aligned arrays of waves transitioning to staggered arrays. c: Smoke tracer concentration.⁷ d: $|\omega_z^+|$. Aligned arrays of waves become chaotic, where chaos spreads from the corners. e: Smoke tracer concentration, transition.⁵ f: $|\omega_z^+|$, uncoupled ($k_x = 0$). $\mu = 16(1 + 0.1i)$, $\lambda = (1 - 2i)/A_0$. IC: Single pulse in ω_z near $z^+, t^+ = 0$; transition. Corner BC. g: Dye tracer concentration of turbulence spot from Prof. M. Gad-el-Hak (private communication).

Note: Figures 1c and 1e are sourced from references 7 and 5, respectively.

3. SHARK DERMAL DENTICLES ARE TEMPLATES OF S-WAVE TRIADS

In a TBL, the patterns in the distributions of the z^+ gradient of ω_z^+ phase, $\partial\phi(\omega_z^+)/\partial z^+$, resemble the irregular distributions of pockets¹⁵ and shark scales (figures 2a – 2c). The contours of ($|\omega_z^+| < \text{threshold } 0.3$) are similar to published high- and low-speed streak patterns.²⁷ The average streak spacing agrees with other published data (figures 2d – 2f and B-2).

Simultaneous observations of high and low frequencies throughout the TBL,^{19,20} and the formation of streaks (solitary waves) and carrier waves due to crashing lateral diffusion waves (figure 4f, videos SI-2,3), can be synthesized in a low-/high-frequency wave-guide model of wall-pressure fluctuations (figures 2g,h). From the converging/diverging dislocation singularities, a pair of streaks and hairpins (they have viscous cores) form a closed laminar viscous resonant “cavity” (strictly speaking, the “cavities” are partially connected) (figure 2g). Analogy is drawn to Schumann resonance in the ionosphere cavity, replacing the speed of light with that of sound, the earth’s circumference with the hairpin vortex tube length, and the ionosphere’s height with the vortex tube diameter. Two disparate groups of resonant pressure waves, large and small, respectively, are produced by the cavity length and diameter as $\omega_n v / U_\tau^2 = [c^+ / (\Lambda^+)] \sqrt{n(n+1)}$, where ω_n is the angular frequency, v is the kinematic viscosity, c^+ is the ratio of the speed of sound in the medium c to the friction velocity U_τ , $\Lambda^+ = \Lambda U_\tau / v$ is the length L^+ or the diameter d^+ of the “cavity”, and n is the mode. Losses and leaks to other cavities would reduce the measured frequencies and widen the spectral peaks, but the resonance would be long lasting, yielding spatio-temporal/TBL pressure fluctuation with two regimes of wall-pressure spectral density (PSD) $\Phi \propto \omega_n^{-N}$, where N is a different constant in the high- and low-frequency ranges of resonance. The wave-guide model explains why there is a clear kink in the PSD and Reynolds number independence (figure 2h). Typically,³² for a streak length of 1000 wall units (figure 2e) and $\delta^+ = 15,800$ at a momentum thickness Reynolds number of 36,000, $L^+ = \frac{LU_\tau}{v} = 47,000$ (considering hairpins inclined at 45° to the flow direction), and $d^+ = 200$ wall units. For $n = 1$ to 3, the model yields $\frac{\omega_{low} v}{U_\tau^2} = 0.0102 - 0.0250 (< 0.5)$, and $\frac{\omega_{high} v}{U_\tau^2} = 2.37 - 5.81 (> 0.5)$. Compared to the time-averaged wall-pressure measurements of PSD, these values correctly lie in two distinct Reynolds-number-independent PSD decay regions, the slope changing abruptly near $\frac{\omega v}{U_\tau^2} = 0.5$. In contrast with prevailing notions, ω_{high} may be sourced from anywhere in the rotational TBL, and ω_{low} solely from the society of interconnected wall-streaks without the hairpins.

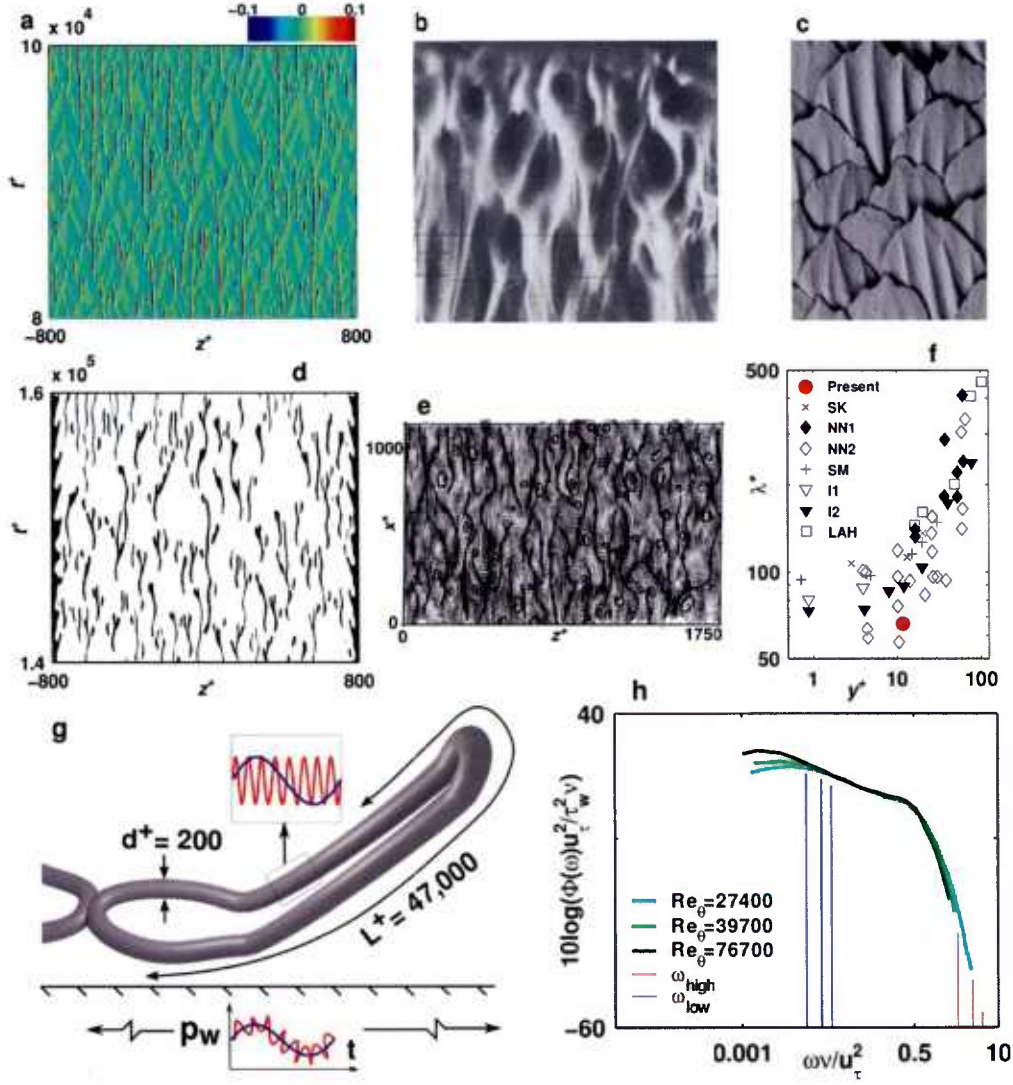


Figure 2. Shark Skin (c) Is a Negative of the TBL Instability Map (a, b)

a: $\partial\phi(\omega_z^+)/\partial z^+$ in (z^+, t^+) ; TBL theory; compare with **b** and **c**. **b:** Near-wall smoke tracer concentration in (z, x) in TBL;¹⁵ $y^+ = 11-13$. **c:** Dermal denticles (diamond-shaped scales) of great white shark;³³ depending on species, there may be three to six riblets per scale (five to six riblets are rarely seen). **d:** Present theory, uncoupled $|\omega_z^+|$ ($k_x = 0$). $\mu = 16(1 + 0.1i)$, $\lambda = (1 - 2i)/A_0$. $|\omega_z^+|(z^+, t^+) < \text{threshold} = 0.3$, which present authors call (ω_z^+) streaks. **e:** Low- and high-speed streak maps obtained using PIV²⁷ in (z^+, x^+) . **f:** Measurements of mean lateral streak spacing²⁷ (λ_z^+, y^+); red circle: present spatio-temporal-averaged value (figure B-2b). **g:** Schematic of origin of wall-pressure fluctuations (p_w) showing vortex-tube wave-guide model of resonance. **h:** Modeled resonance frequencies (ω_{high} , ω_{low}) compared with measurements³² of spectra ($\phi(\omega)$) in TBL of high momentum thickness Reynolds number (Re_θ); three vertical lines in each of the ω_{high} and ω_{low} ranges, represent (from left to right) modes 1, 2, and 3, respectively.

Note: Figures 2b, 2c, and 2e-f are sourced from references 15, 33, and 27, respectively.

Figure 2h is adapted from reference 32.

4. TBL IS PERENNIALY AND CHAOTICALLY TRANSITIONAL

The thickness of the molecular viscosity-dependent wall layer undergoes limit cycle oscillation between approximately 5 and 50 wall units (figure B-1d). It blends with the viscous-sublayer near the wall and temporal inner extrapolations of the log-layer at the other edge. The phase-averaged streamwise velocity profile in the y^+ direction approaches a modified Blasius relationship as the sublayer thickness approaches $y^+ = 50$ near bifurcation (figure B-1). The appearance of TS (Tollmien-Schlichting) wave-like, short-span, two-dimensional oscillations in the wall layer is expected in the pre-bifurcation region in this latter part of the sublayer growth (see figure B-1 and the Supplementary Information (SI) in appendix C); these oscillations are termed pseudo-TS waves (S TS) (see Methods). The distortions of these waves are shown in figures 3a-l, which present iso-contours of $\phi(\omega_z^+)$, as a series of interrupted parallel lines.

Streaks are formed where a temporal series of diffusion waves ends (figures 3a, 4b). They are straight in arrayed transition (figure 3a, left). In a TBL, the streaks meander (figure 3a, right) due to unequal lateral phase-shock velocities³¹ approaching from the opposite boundaries (figures 4d, 4f).

Figure 3c (bottom) shows synchronous oscillations. Large and small phase mismatches in successive oscillations^{21,22,31} distort the phase maps. In figure 3b (top and bottom), tongues form due to phase mismatch between prior and later oscillations. Figure 3c (top and middle) shows oblique waves. The triad base (black arrow) is the hairpin nucleation site (figure 4c). Figure 3d shows numerical results³⁴ of staggered iso-shear contours in H-type transition for comparison.

In figure 3e (top), waves end; in figure 3e (bottom), waves merge and come to rest when they have a large phase difference. Since the instabilities originate in two-dimensional waves, the number of waves ending and merging is equal.

Within the transitional range that exists before chaos ensues, the waveforms are of K- and H-types⁷ (figure 3f); these are, respectively, *fairly precisely* aligned and staggered waveforms. After chaos ensues in these waveforms, the flow is termed a TBL. Here one also finds two kinds of instabilities, aligned and staggered,³⁵ but *approximately* so, where the latter grows spatially (subharmonic formation as in H-type) but the former doesn't. These instabilities are named A-type for aligned waves and S-type for staggered waves. Also, they are called triads since they have two oblique waves and a two-dimensional oscillation at the base (the former is formed when, locally, the vorticity fluctuates differently from its neighbors—a so-called pacemaker). The S-type wave gives rise to nucleation sites of hairpin vortices at its base (figures 4a and 4b).

Figures 3f and 3g point out (confirmed later in figures 5 and 6) that *in the journey from H to S, the arrest of the growth of the S-triads is synonymous with the preservation of H-type transition and chaos control*. In figure 3h, the seemingly complicated pattern is actually composed of groups of interacting S- and A-triads.

The corner and periodic boundary conditions produce variations of instabilities. The former produces a Strouhal shedding of hairpin vortices (blue at right in figure 3j and blue outlines at left in figure 4b), with oscillating interaction with oblique waves (figures 3i, right; 4b; B-2a), while periodic boundary conditions produce pure oblique waves (figure 3i, left). In figure 3i (right), in corners, waves end and merge.

In a TBL, there is sometimes an avalanche of interactions between the S- and A-triads in which the S-type jumps streaks, laterally swallowing an array of A-triads (figures 3j and 3k). The process can continue downstream, and repeated subharmonic growth fills the spectrum; note that in contrast to H-type, S-waves can grow. Apart from these global patterns, locally the vorticity waves can end or merge (figure 3e) or oscillate (arrow in figure 3c) between them. Interactions of both wave triads yield spatio-temporal self-regulation (figure 3k).

In figure 3k, the group patterns are unstable because the antagonistic diffusion is slow. In biology, stable patterns are produced in embryos when the head-to-tail and left-to-right axes are being determined (with some handedness), the signal for this vanishing soon after. Similarly, K- and H-types of transition are produced when the IC and BC are felt, and then K- and H-types become unstable (the effect of handedness growing to further bifurcation), which is called a TBL. Therefore, the activator (a) and inhibitor (b) equations²¹ in a TBL have many variations of the same embryonic theme, depending on how many bifurcations have taken place (i.e., how high the Re number is). To control chaos, sharks and dolphins manage these variations using competing instabilities.

In figure 3l, at $z^+ = -200$ the vorticity diffusion rates in the activator ($:: a$) and inhibitor ($:: b$) parts of the self-regulation process in S-waves are equal, yielding the boxed “trapped spiral.”

When weakening (less red), the arrayed K-type transitional waves become A-type in a TBL, whose start is given by the onset of chaos (see figure 5k). However, when strengthening, both K- and H-type transitional waves become S-type in a TBL (figure 5f). This becomes clear during chaos control where the instability process is slowed down.

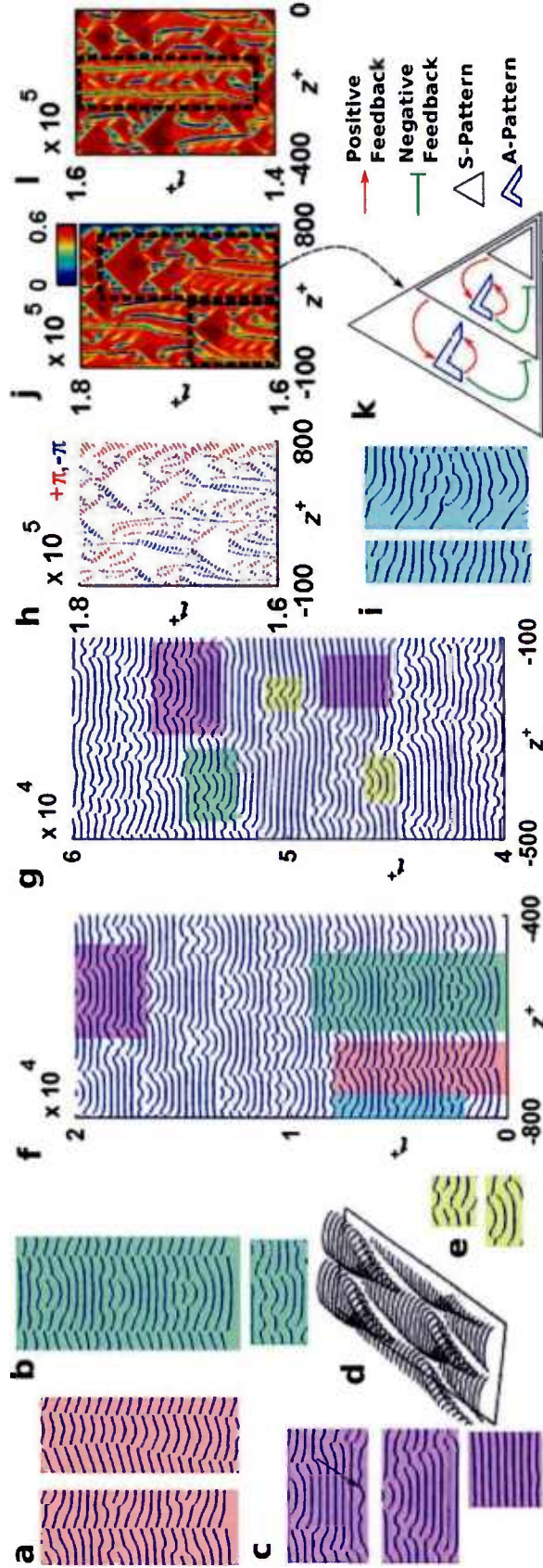


Figure 3. Spatio-Temporal Instability Phase Maps of a TBL

a-c, e-g, i: $\phi(\omega_z^+)(z^+, t^+) = 0.95\pi$; **d:** Numerical simulation of iso-shear;³⁴ **h:** $|\partial\phi(\omega_z^+)/\partial z|/(z^+, t^+) \approx \pi$, $\mu_r = 64$; **j:** $|\omega_z^+|(z^+, t^+)$, $\mu_r = 64$; **l:** $|\omega_z^+|(z^+, t^+)$. **a:** Streaks (left: transition), (right: TBL). **b** (top): K-type transition; **c** (top) and **d:** H-type transition. **e:** wave ending (top), merging (bottom). **f:** K-type to H-type transition. **g, b** (bottom) and **c** (middle, bottom): S-triads. **h, j, l:** TBL. **i:** BC (periodic: left) and (wall: right). Boxes in **j:** Handed avalanche of self-regulating wave interactions. **k:** Schematic of “Pac-man” model of self-regulation of wave interaction; red arrow: autocatalytic activator *a*, blocked green arrow represents *b*, *a*’s antagonist, the inhibitor.²¹ **l:** Box marks S-triads oscillating between streaks, giving rise to a trapped spiral because the avalanche of streak jumps (**j, k**) is not triggered.

Note: Figure 3d is sourced from reference 34.

5. LATERALLY CRASHING DIFFUSION WAVES PRODUCE WALL STREAKS

Figure 4a shows the distribution of diffusion. In figure 4b, the phase (0.95π) and amplitude ($A_0/2$) of ω_z^+ are overlaid. They both detect the streaks where the iso-phase ends alternately. Hairpin vortices are shed from the left corner boundary in 4b. Figure 4c reproduces the orthogonal plane smoke flow visualization showing the formation of “pockets”¹⁵ and hairpins. Section A-A and inset C₁ show that pockets¹⁵ (figure 2b), folds,⁵ and S-triads are identical waveforms. A hairpin vortex (4c, lower sketch and photograph) forms at the base of the S-triad. After comparison, the hairpin vortex nucleation sites are marked in figures 4a and 4b (markings also in figure 3c (top) and figure B-2a).

Figures 4d – 4f, B-3, and video SI-1 show the simultaneous formation of streaks and small-wavelength oscillations (“carrier waves”) of the state. ω_z^+ streaks form where $Re(\omega_z^+)(z^+, t^+) = 0$. An antagonistic pair of diffusion waves approaches each other laterally.³¹ As they slow down exponentially, nearly coming to a halt, they abruptly devolve into small-wavelength carrier waves. The lateral waves alternately exceed each other’s speed, slightly overshooting the dividing plane and producing a sliding oscillation in t^+ (figures 3a and 4f). (It may be that animals are using inter-aural time differencing of a lateral pair of surface pressure sensors to detect the streaks.) The overshoot causes streak-meandering in a TBL, which is absent in K-transition (figure 3a). To control chaos, periodic surface templates prescribe the overshoot. In figures 4d – 4f, the wave-crashing is depicted in relation to lateral arrays of slightly tilted, near-wall (mushroom) vortex pairs (4d) as a nonlinear pendulum, where the saddle “s”, a point of neutral equilibrium, oscillates.

The formation of carrier waves and lateral crashing are efficient when the IC wavelength of ω_z^+ is $200z^+$. Crashing is not clear when the ω_z^+ wavelength is reduced from 200 (see video SI-2) to 100 (see video SI-3).

The long streaks are like solitary waves with short-wavelength carrier waves riding on them, giving the notion of coupled high and low frequencies.¹⁸⁻²⁰ Since a TS wave is the natural oscillation of the pre-breakdown ν -based diffused near-wall vorticity, the carrier wave is a form of orthogonal ^STS wave. As was shown in section 3, a wave-guide model indicates that the lengths of the viscous cores of streaks (figures 2c, 2e, “f” in figure 4e), which bounce around the wall-parallel TBL, produce low-frequency pressure pulses and their diameters produce distinctly high-frequency pulses.

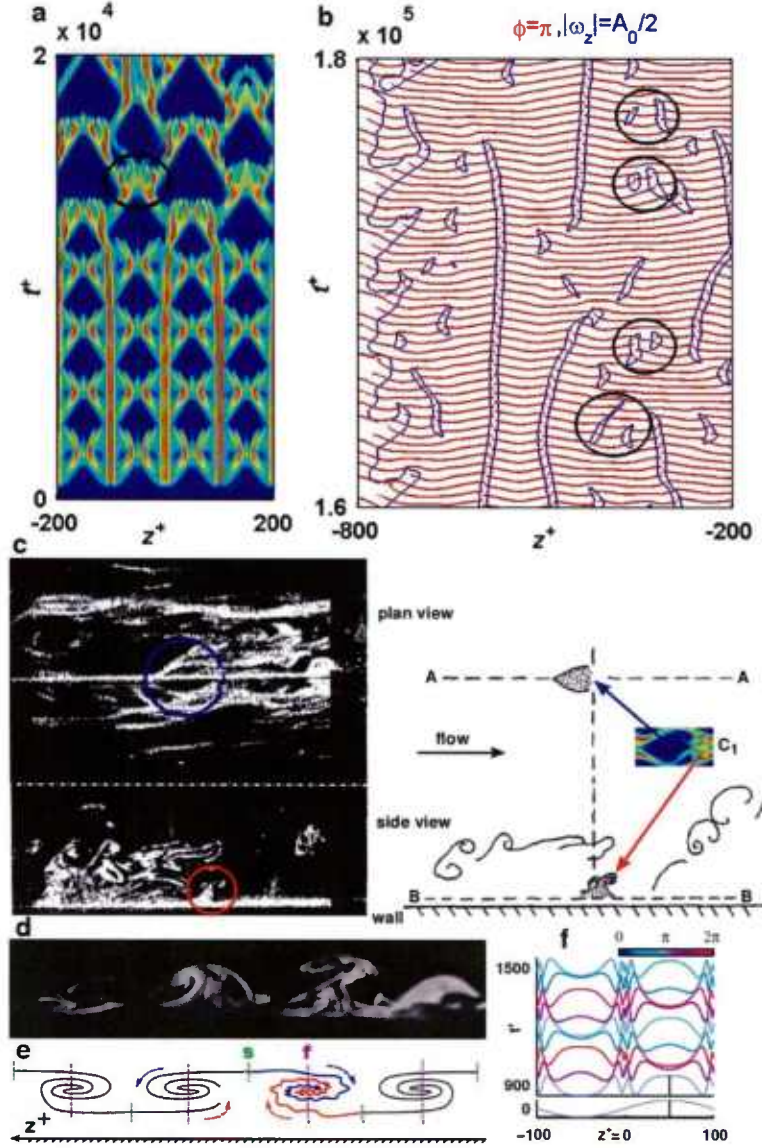


Figure 4. Nucleation Site of Hairpin Vortex and the Origin of Streaks and Small-Wavelength Oscillations

a: $|\partial^2 \omega_z^+ / \partial z^2|$. **b:** Overlay of iso-phase (0.95π) and amplitude ($A_0/2$) of ω_z^+ . **a:** transition; **b-f:** TBL. **c:** Figure 5 of Falco¹⁵ showing simultaneous smoke visualization of surface parallel (top) ($y^+ = 9 - 14$) and longitudinal plane (bottom) TBL. Inset **C₁**: S-triad from **a**. Circles in **a** and **b** are nucleation sites of hairpins. **(d-f):** Simultaneous formation of streaks and small-wavelength oscillations of the state. **d:** Cross-stream (z, y) near-wall smoke tracer concentration;⁵ the mushroom-like longitudinal vortex pairs are slanted. **e:** Schematic of nonlinear pendulum behavior of vorticity diffusion; phase: (**s**: saddle, unstable 180°) and foci (**f**: stable, 0°). **f** (Upper Traces): $\text{Re}(\omega_z^+)(z^+, t^+)$, $\mu_- = 0.5$. ω_z^+ scale bar at $z^+ = 50$: $2\sqrt{\left(\frac{\sigma_r}{\lambda_r}\right)} = 1.2$. Note crashing of lateral diffusion waves and formation of carrier waves. Color indicates $\phi(\omega_z^+)$ at $z^+ = 50$. **f** (Lower Trace): $\text{Re}(\omega_z^+)(z^+, t^+ = 0)$; scale bar at $z^+ = 50$: 0.2.

Note: Figures 4c and 4d are sourced from references 15 and 5, respectively.

6. SHARK SKIN IS A SPATIO-TEMPORAL DIFFUSION TEMPLATE FOR RESETTING THE PHASE OF CHAOTIC WAVE TRIADS

Figure 5 shows the results of the modeling of chaos control by sharks. Skins of great white,³³ Atlantic sharpnose,³⁶ and tiger shark³⁷ dermal denticles (dd), with embedded riblets, are shown in figures 2c, 5a, and 5e, respectively. The model using a bi-modal periodic distribution (in z , t , or z , x) of the coefficient of lateral diffusion μ (see Methods) is shown in figures 5b and 5d. The agreement of figures 5b and 5d with the $|\omega_z^+|$ of TBL S-waves (figures 5c and 5f) and shark skins (figures 5a and 5e) is remarkable.

Figure 5f shows the baseline TBL $|\omega_z^+|$. Figure 5g is a diffusion contour map of the riblets (R), the dermal denticles (dd), and of their combination (R+dd). At the uninterrupted riblet spacing of $10z^+$, lateral diffusion (k_μ) perturbations of 2.5% (weak riblets) delay the onset of chaos, and K-type waves are conserved for $t^+ \leq 80,000$ (figure 5h); a 20% perturbation (strong riblets) then applied at $t^+ \geq 80,000$, the chaotic region in 5h, produces orderly arrays of A-waves, suppressing S-waves (figure 5i). However, the riblets need to be interrupted along the dd outlines (5d,e,g) because the animal is flexible and non-uniform in x .

Figure 5j shows the combined effects of weak riblets that are interrupted by the dd (see Methods). Figure 5k shows the effects of combining weak interrupted riblets, dd, and strong diffusion perturbations after chaos has ensued ($t^+ > 80,000$). This is the predicted chaos control due to sharks.

See figures B-4 through B-7 and the SI for additional results.

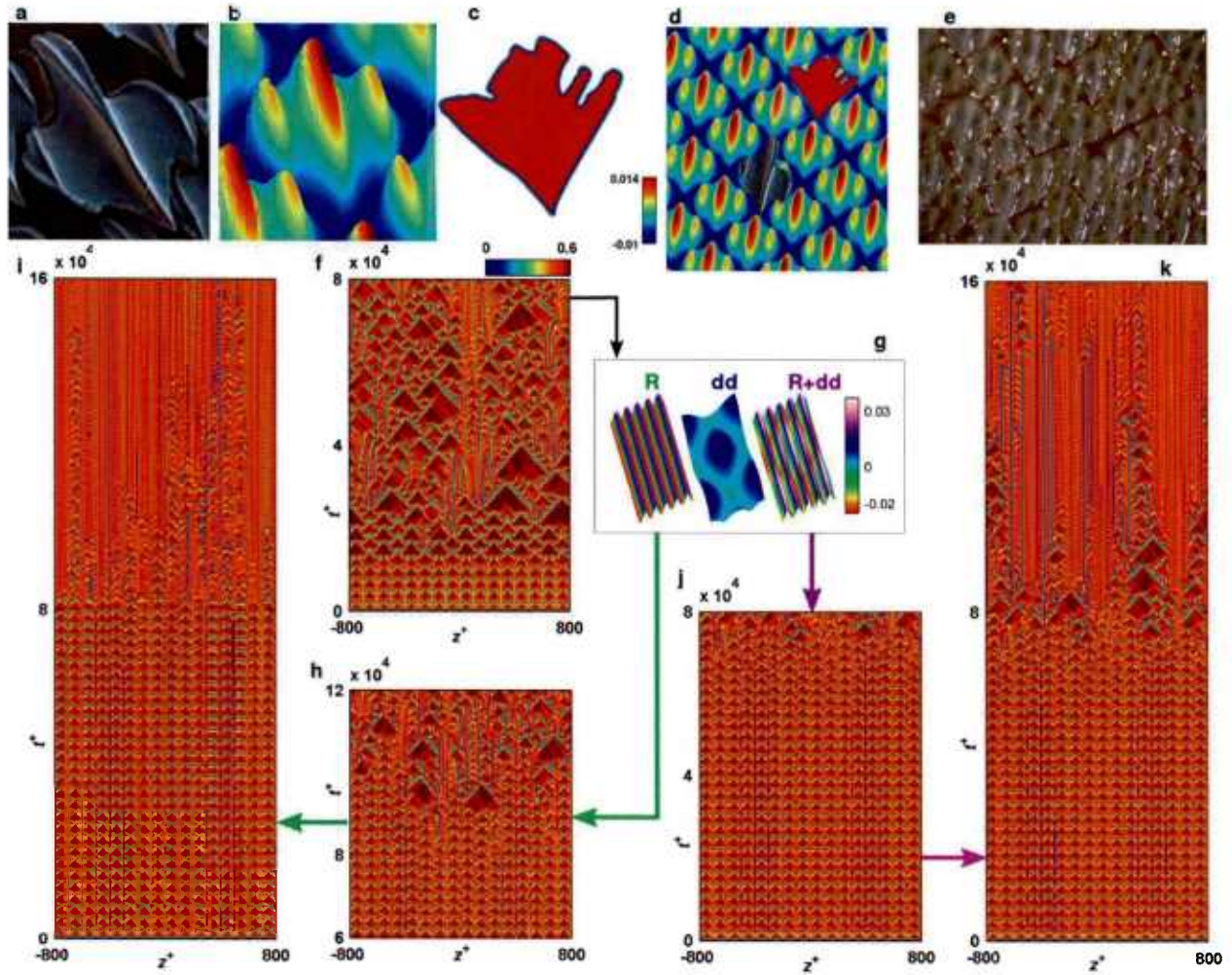


Figure 5. Shark Skin Is a Template of Organized Lateral Diffusion for Control of Chaos Onset

a: Atlantic sharpnose dd.³⁶ **b:** Diffusion ($\mu(z^+, t^+)$) analog of **a**. **c:** $|\omega_z^+|(z^+, t^+)$. A single S-triad is shown. **d:** Modules **a** and **c** are overlaid on a staggered model of **b**. **e:** Staggered array of dd in tiger shark³⁷ compares with **d**. **f-k:** $|\omega_z^+|$ maps; riblet spacing $s^+ = 10$. **f:** Baseline TBL. **g:** μ -analog of surface: **R**, **dd**, and **R+dd** refer to riblets, dermal denticles, and combined riblets and dermal denticles, respectively. **h:** Uninterrupted weak riblets; chaos onset is delayed; arrayed K-waves are conserved. **i:** Uninterrupted strong riblets at $t^+ \geq 80,000$; arrayed A-waves are produced. **j:** Combined weak riblets of smoothly varying heights as in **dd**. **k:** Weak, then strong riblets and **dd**.

Note: Figure 5a, overlay in figure 5d, and figure 5e are sourced from references 36, 36, and 37, respectively.

7. DOLPHIN LATERAL MICROGROOVES VIBRATE TO RESET PHASE OF CHAOTIC WAVE TRIADS

Figure 6 shows how dolphins control chaos with surface perturbations coupled to the SL oscillator. The baseline TBL in figure 5f applies. In dolphins, the lateral surface microgrooves (inset A in figure 6)³⁸ organize the lateral coupling and a two-dimensional SL oscillator is approached (see Methods); the ^STS waves are prominent in figures 3c (bottom) and 4b.

The dolphin's microgrooves impart variation in the coefficient of lateral diffusion whose periodicity matches that of ^STS waves. Active vibration³⁹ of the microgrooves, modeled as I_{ext} (see Methods), is assumed to be in integer multiples (1, 16) of ^STS waves (figures 6a and 6b). With bimodal vibration (1^STS + 16^STS), chaos onset is delayed the most (figure 6c). With grooves alone, fewer S- but more A-triads are produced (figure 6d). When microgrooves and vibrations (16^STS) are combined, there is a clear delay in the onset of chaos (figure 6e).

The diffusion template lock-in mechanism of chaos control by sharks and dolphins is a spatio-temporal version of Llinas' olivo-cerebellar temporal SPR mechanism¹ (figure B-10). In the latter, the relative physical locations of the inferior-olive neurons are not important for phase synchronization, while the spatial locations of the oscillators are important in TBL chaos control.

Riblets in sharks are longitudinally proud and periodic in z^+ . Dolphins can have longitudinal arrays of transverse microgrooves in x^+ (inset A in figure 6). Both surface patterns control diffusion. Compared with the baseline in figure 5f, the S-triads in figure 6e are fewer and there is a corresponding increase in A-triads; the $|\omega_z^+|$ map is similar in riblets (figure 5i).

When chaos is controlled, turbulence production-induced drag is reduced, the limit being 50%.³⁰ The present authors therefore see nothing in their theory that contradicts the prior claims^{40,41} regarding dolphins' high drag reduction.

Engineering drag reduction methods, such as outer layer devices⁴² (figure B-8) and Stokes^{30,43} spanwise shaking of walls (figure B-9), potentially can be modeled as periodic perturbations that are *external* (I_{ext}) to the SL oscillator (lacking in coupling).

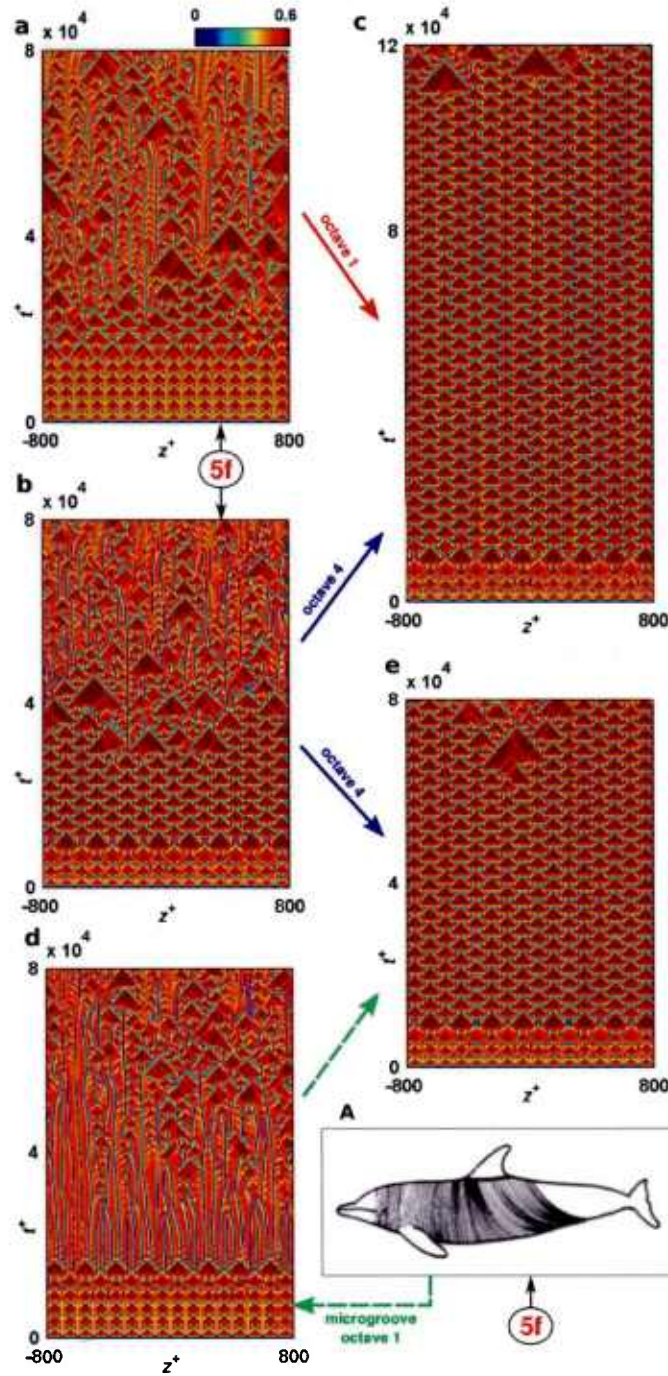


Figure 6. *Dolphin Chaos Control with Surface-Grooved Diffusion Perturbations Coupled to the SL Oscillator, and External S TS-Wave Microvibration of Surface*

a – e: $|\omega_z^+|$ maps. Baseline TBL is as in figure 5f. Surface microvibration: 1^S TS (a); 16^S TS (b); 1^S TS + 16^S TS (c). **Inset A:** Microgrooved dolphin.³⁸ **d:** Effect of microgrooves. **e:** Combined lateral microgrooves and 16^S TS vibration.

Note: Inset A in figure 6 is sourced from reference 38.

8. DISCUSSION AND SUMMARY

Wall turbulence models, on which some climate models may be based, tend to be statistical, although there is no instant when any of the states have profiles matching these models. Economics theories have similar artificial reality. Newton's laws of motion, on the other hand, apply at all instants, and the present theory does as well. If IC and BC are sufficiently known, the present theory makes a short-term theoretical prediction. It describes TBLs using orthogonal SL oscillators diffusively coupled laterally. Sharks and dolphins control the growth of chaos in turbulence and maintain the existing transitional organization by (1) imposing a weak lateral and longitudinal periodic perturbation to the oscillator's diffusion matching their skin topology, and (2) imposing periodic perturbations external to the oscillator much as in Llinas's¹ olivo-cerebellar synchronization, again using their skin.

Turbulence production over a wall is highest near $y^+ = 10 - 15$. The theory considered instabilities in a layer parallel to the wall at this height. Due to wall proximity, the coupling with the third orthogonal oscillator (ω_y^+) was ignored. In the unstable wall-parallel layer at $y^+ \approx 11$, unsteadiness of the pre-breakdown laminar flow is modeled by a parabolic equation⁴⁴ (equation A-2 in Methods), where reversal of the direction of vorticity diffusion makes the equations singular parabolic, whose attribute is irregularities within the flow field. In the present case, this gives rise to rows of stagnation points (breaks in ω_z^+ contours—figure 3a) appearing as streaks, and dislocation sites where streaks merge or diverge. The boundaries where diffusion changes sign are regions of large variation in tracer concentration and density (figures 3b and 3e), and where chaos ensues (e.g., at the K to H transition, as slightly-handed streak jumps). The surface templates in sharks/dolphins control the motion of the stagnation points by ensuring similar handedness in bifurcations²⁴ of the diffusion wavefronts, restricting the diversification of scales.

The Reynolds numbers (Re) of whales (up to blue whales) ($0.4-3 \times 10^9$) and dolphins ($0.2-1.75 \times 10^8$) are similar to those of ships/submarines ($> 2 \times 10^9$) and unmanned underwater vehicles ($> 2 \times 10^7$), respectively. The feasibility of cost-effective chaos control in practical flows is not in doubt. The removal of I_{ext} establishes a new TBL having higher lateral mixing (figure B-8b), and dolphins may also delay the onset of chaos up to almost halfway between the dorsal fin and the tail (inset A in figure 6), beyond which chaos' enhanced mixing keeps the TBL attached. It would be useful to document the dynamic skin properties of whales whose high Re number may involve chaos management mechanisms similar to high-frequency neural control.⁴⁵

Currently in weather simulation, the mean paths of hurricanes are computed but no predictions are offered for tornadoes due to spatial resolution issues. Alternatively, since a self-regulating model of weather should be valid,^{46,47} a low spatial resolution model would produce accurate results if weak orthogonal oscillators are included as control oscillators. Four-equation riblet modeling carried out with low spatial resolution (see figures B-7a, c, e) reproduces the measurements qualitatively. Tuning such oscillators with high spatial resolution is difficult; high-resolution simulations with fewer orthogonal coupled oscillators have been emphasized here instead. Even a one-oscillator model (ω_z , figures 1f, 2d, B-4, B-7f-h, and B-8b,c) captures some of the TBL spatio-temporal organization. Persistent chaos control, as animals must be doing, requires inclusion of all three orthogonal coupled vorticity oscillators of amplitude 1,

0.01, and 0.0001 (in the z , x , and y directions, respectively), and of the viscous layer thickness, which has a mass conservation property.

For the delay of chaos onset, the perturbations required are weak. But the diverse interactions of S- and A-wave triads produce a Reynolds number-like effect (see figure 7D₂ in Kazantsev et al.¹), which requires larger perturbations for control.

Laterally growing non-uniformities and end-effects are present even in laminar boundary layers (LBLs). LBLs approaching Re_{crit} become unsteady and can be described by highly diffusive SL equations with periods matching those of TS waves (see Methods). Such two-dimensional waves are seen in TBLs, and during control they become clear (figure B-8b). During swimming, dolphins also develop arrays of parallel azimuthal microgrooves not only near their nose, but over most of their length (see inset A in figure 6).

The lateral crashing of diffusion waves in a TBL is similar to oblique shedding in cylinders spreading from the ends to the center at the speed of phase “shock.”³¹ While the crashing of diffusion waves and the formation of carrier waves and streaks remain qualitatively unchanged with lateral diffusion, the carrier waves are clearer at low values ($\mu_r = 0.5$) (figures 4f and B-3). When μ_r drops from 64 to 0.5 in $t^+ = 1700$, the number of crashes increases from four to five, which is a Reynolds number-like effect. Although results are given here primarily for $\mu_r = 16$, μ_r should change in steps as the near-wall turbulence instability goes through bifurcations to produce ^STS waves; arrayed, staggered, and chaotic triads; and their more intricate interactions.

The present approach is compatible with common engineering drag reduction methods, which are modeled as perturbations external (I_{ext}) to the TBL oscillators (figures B-8 and B-9).

Lateral diffusion of vorticity forms streaks and small-wavelength carrier waves, but evidence of their coupling is lacking. A society of laterally balancing streaks form *ad infinitum* on a whale from head to tail—but the size of the largest scales is limited by the competition of the activator and inhibitor parts of the self-regulation process of instabilities.

Only the phases of non-uniformities that are lower or higher even-harmonics of ^STS wavelength are reset to null the instigation of chaos because other wavelengths do not couple with the TBL oscillators. This accrues from the disturbance rejection property of SL equations. This property limits the work cost of control, which would be less than previously thought.

Dolphins use wavelengths of (^STS, 16^STS) for longitudinal surface vibration acting on the ω_z^+ oscillator for control. Dolphins’ fatty skin is conducive for this, whereas sharks’ dermal denticles and riblets are made of rigid teeth-like material that acts like armor. In sharks (figures 5f-5k), external perturbation did not yield large benefits. In the evolutionary reaction-diffusion bifurcations,^{21,22} sharks and dolphins have taken more aggressive and docile turns, respectively, and their chaos control strategies are consistent with that selection. The hydrodynamics, control, and sensing of sharks/dolphins’ flapping fin propulsion are self-regulating,²⁴ just as their wall-turbulence is. Both flapping fin hydrodynamics and wall-turbulence are transitional,³⁵ and have

bi-stable attractors. The mechanism of how these animals use their actuators, sensors, and controllers to interact with their surroundings has a common theoretical foundation.

The self-regulating, transitional, and deterministic nature of a TBL at all Re numbers makes chaos control an exercise in finding the locally effective spatio-temporal phase reset mechanism. The perturbations of tornadoes in planetary boundary layers are likely to have predictable control solutions, just as sharks and dolphins likely deal with the oscillating fin-necklace vortices submerged in their TBLs. Theoretical insights into cerebral electrical storms during seizures could be gleaned by modeling neuronal networks as dislocating three-dimensional streaks.

This report has shown chaos control by surface “emplating” (evolutionary interactive templating between the turbulence environment and the animal skin) and external stimulation. In a broadened definition, nascent chaos control may be deemed to be present in other systems in nature, such as economics regulation, parental guidance, social bonding, rhyming naming of siblings, and large-scale societal regimentation.

APPENDIX A METHODS

METHODS SUMMARY

Evolution of spanwise and streamwise vorticity disturbances in a TBL is modeled by two diffusively-coupled continua of Stuart-Landau oscillators, derived near the critical Reynolds number from the Navier-Stokes equations. These oscillators are orthogonally coupled to each other in a manner consistent with continuity. The equations apply in the pre-breakdown TBL, a region of time-variable thickness dominated by molecular viscosity rather than eddy viscosity. The equations' parameters are derived from published experimental values measured in the post-transition ($Re_x \gg Re_{x,crit}$) TBL. TBL control is accomplished by spatio-temporal variation of the diffusion and/or by an external source of vorticity. The details are given below.

APPLICABILITY OF THE STUART-LANDAU EQUATION

From the Navier-Stokes equations, Stuart⁴⁸ and Watson⁴⁹ showed that small disturbances $A(t)$ in parallel (Couette, Poiseuille) flows near the critical Reynolds number Re_{crit} evolve via the Stuart-Landau (SL) oscillator equation (A-1). This model is broadly applicable²⁸ for the growth of disturbances near a bifurcation point:

$$\dot{A} = \sigma A - \lambda A^2 A^*, \quad (\text{A-1})$$

where A is any physical quantity indicated by the Navier-Stokes equations,³¹ and $(\dot{\cdot})$, $(\cdot)^*$ indicate time derivative and complex conjugate operations, respectively. The complex coefficients σ , λ are determined by the flow and modeled quantity. A forced form of equation (A-1) can quantitatively model the force fluctuations on a flapping fin.²⁴ The model is applicable to non-parallel flows if the flow development time scale is slow compared to that of perturbation growth.

Spanwise diffusive coupling can be introduced into equation (A-1), such that

$$\dot{A}(z, t) = \sigma A + \mu \frac{\partial^2 A}{\partial z^2} - \lambda A^2 A^* \quad (\text{A-2})$$

Equation (A-2) was proposed⁵⁰ for a bluff body shedding model composed of discrete Van der Pol elements, which are a special case of SL oscillators.³¹ This coupling is the viscous force produced by Newtonian fluid between closely-spaced oscillators, and approaches the differential form as spacing goes to zero. This model reproduces the growth of chevron patterns in the cylinder wake.³¹

Einstein and Li²⁶ observed a time-dependent viscous layer in the near-wall region of the TBL. This layer's growth rate is inversely proportional to its thickness and is liquidated when its thickness renders it unstable. This liquidation ejects the built-up vorticity into the outer flow,

and then a new cycle resumes. This sublayer's existence is required to employ the SL equation in the TBL because the assumption $Re \approx Re_{crit}$ required to employ the SL model is valid there.

NEAR-WALL TBL MODEL

Spanwise and streamwise vorticity perturbations, ω_z^+ and ω_x^+ , respectively, were modeled by the SL equations

$$\dot{\omega}_z^+ = \left(\sigma \omega_z^+ + \mu \frac{\partial^2 \omega_z^+}{\partial z^{+2}} - \lambda^z \omega_z^{+2} \omega_z^{+*} - k_x |\omega_x^+| \right) \tau, \quad (\text{A-3})$$

$$\dot{\omega}_x^+ = \left(\sigma \omega_x^+ + \mu \frac{\partial^2 \omega_x^+}{\partial z^{+2}} - \lambda^x \omega_x^{+2} \omega_x^{+*} + k_z \frac{\partial \omega_z^+}{\partial z^+} \right) \tau, \quad (\text{A-4})$$

where τ is a time scaling and the parameters k_x and k_z are positive, real coupling constants. The notation $(\cdot)^+$ indicates a quantity in wall units. The coupling mechanism indicates a sink in ω_z^+ scaling with ω_x^+ , and a source in ω_x^+ consistent with continuity, assuming that $\omega_y^+ \ll \omega_x^+$ and $\partial/\partial x \approx \kappa \partial/\partial t$. This is called an ‘‘orthogonal’’ coupling due to this consistency with continuity. The simplest mechanism meeting these requirements was chosen, but others exist.

PARAMETER ESTIMATION

The SL oscillator in equation (A-3) has base amplitude $A_0 = (\sigma_r/\lambda_r^z)^{1/2}$ and spanwise scale proportional to $(\mu_r/\sigma_r)^{1/2}$, where $(\cdot)_r$, $(\cdot)_i$ indicate the real and imaginary parts, respectively, of a complex quantity. Near the wall, the wall-shear stress $\tau_{wall} \approx \omega_z$, since $dv/dx \approx 0$ (v is the surface-normal velocity). The quantity $\tilde{\tau}_{wall}/\bar{\tau}_{wall}$ is in the range⁵¹ $0.32 - 0.40$, where $\widetilde{(\cdot)}$ and $\overline{(\cdot)}$ are the standard deviation and mean, respectively. In the present model, $A_0 = 0.63$ yields $\tilde{\omega}_z^+ = 0.364$; $\bar{\omega}_z^+ = 1$ by definition. For $\sigma = \sigma_r(1 + c_0 i)$, $\mu = \mu_r(1 + c_1 i)$, and $\lambda^z = (1 + c_2 i)/A_0^2$, $c_1 = 0.3$, $c_2 = -3$ have been employed (unless noted). $c_0 = 1$ yields unit linear frequency.

Unless noted, the spanwise scaling factor $\mu_r = 16$. This is estimated from³¹ $\sigma_r - \mu_r(\pi/L_c)^2 = s(Re - Re_{crit}) = 0$ for characteristic aspect ratio L_c . $Re = Re_{crit}$ is assumed per the self-regulating nature of the sublayer. L_c is the ratio of near-wall streak spacing to measurement layer height (figure 2f²⁷). The best fit for points in the range $0 \leq y^+ \leq 50$ is $\mu_r = 467.7(y^+)^{-1.62}$. Using this fit, $\mu_r = 16$ corresponds to $y^+ = 8.10$ near the range $9 \leq y^+ \leq 14$ indicated by comparison of the present results with spatio-temporal observations.¹⁵

Townsend⁸ derived $\tilde{\tau}_{wall,x}/\tilde{\tau}_{wall,z} = 1/100$ using assumptions similar to those used in the present work. It is then estimated that $\lambda_r^x = 100^2 \lambda_r^z$. Order-of-magnitude considerations yield $k_x = 1$ and $k_z = 0.01$. $\tau = 0.005$ was chosen for computational convenience, since the oscillator time scale is arbitrary.

PERIOD OF A TWO-DIMENSIONAL UNCOUPLED SL OSCILLATOR

Let $\omega_z^+(z^+, t^+) = M(z^+, t^+) \exp(i\phi(z^+, t^+))$, where magnitude $M \in \mathbb{R}$, $M \geq 0$, and phase $\phi \in \mathbb{R}$. In regions where $M(z^+, t^+) \approx M(t^+)$, $\phi(z^+, t^+) \approx \phi(t^+)$, equation (A-3) simplifies to

$$\dot{M} = (\sigma_r M - \lambda_r^z M^3) \tau, \quad \dot{\phi} = (\sigma_i - \lambda_i^z M^2) \tau,$$

such that for $M > 0$, $\lim_{t^+ \rightarrow \infty} M = A_0 = (\sigma_r / \lambda_r^z)^{1/2}$. An analogous limiting value $\dot{\phi}_\infty = (\sigma_i - (\sigma_r / \lambda_r^z) \lambda_i^z) \tau$ yields the period of the two-dimensional (Tollmien-Schlichting) wave $T = 2\pi / \dot{\phi}_\infty$, where $T = 100\pi$ for the constants used here. This period is used in the modeling of microgrooves and microvibrations in dolphin chaos control.

SIMULATIONS

The equations are solved with a finite-difference fourth-order Runge-Kutta solver with time step $dt^+ = 0.2$ and spatial grid $dz^+ = 1$. Harmonic boundary conditions are applied at $z = \pm L/2$ unless noted. When used, corner boundary conditions $\omega_z^+(\pm L/2, t^+)$ are applied in combination with $\sigma_r = \sigma_r(z^+)$ near the boundary; to obtain $\sigma_r(|z^+| - L/2) \approx \sigma_{r,center}$ and smooth descent to zero at the boundary, $\sigma_r(z^+) = \sigma_{r,center} \times [\text{erf}((z^+ + L/2)/25) - \text{erf}((z^+ - L/2)/25) - 1]$. The value $\sigma_{r,center} (= 1)$ is the σ_r value used for all z when using harmonic boundary conditions. This equation yields a “corner” width of $z^+ \approx 50$, matching Re_{crit} . Unless noted, the initial conditions, which embody lateral cellularization due to end-effects, are $\omega_z^+(z^+, 0) = 0.1 \sin(2\pi t^+/200)$, $\omega_x^+(z^+, 0) = 10^{-5} N(z^+)$, where $N(z^+)$ are uniformly distributed random numbers with range $(-0.5, 0.5)$ (N simulates free-stream turbulence).

CONTROL

“Controlled” forms of equations (A-3) and (A-4) are:

$$\dot{\omega}_z^+ = \left(\sigma \omega_z^+ + \mu^z(z^+, t^+) \frac{\partial^2 \omega_z^+}{\partial z^{+2}} - \lambda^z \omega_z^{+2} \omega_z^{+*} - k_x |\omega_x^+| + I_{ext}^z(z^+, t^+) \right) \tau,$$

$$\dot{\omega}_x^+ = \left(\sigma \omega_x^+ + \mu^x(z^+, t^+) \frac{\partial^2 \omega_x^+}{\partial z^{+2}} - \lambda^x \omega_x^{+2} \omega_x^{+*} + k_z \frac{\partial \omega_z^+}{\partial z^+} + I_{ext}^x(z^+, t^+) \right) \tau.$$

The physical mechanism of “ μ -control” is alteration of the local aspect ratio L_c . The mechanism of “ I_{ext} -control” is the addition of vorticity from a source outside the model.¹

SHARK CHAOS CONTROL

Uninterrupted riblets are modeled using $\mu^z(z^+) = \mu^x(z^+) = \bar{\mu}(1 + \kappa_\mu \sin(2\pi z^+/s^+))$, where $\bar{\mu} = 16(1 - 0.3i)$ is the flat plate value. The spacing between riblet peaks $s^+ = 10$ is used unless noted. κ_μ gives the “strength” of the riblets, where the terms “weak” and “strong” riblets denote $\kappa_\mu = 0.025$ and 0.2 , respectively.

Riblets interrupted by the dermal denticles are modeled using $\mu^z(z^+, t^+) = \mu^x(z^+, t^+) = \bar{\mu} (1 + \kappa_\mu \sin(2\pi z^+/s^+) + Sc(z^+, t^+) - \overline{Sc})$, where $Sc(z^+, t^+) = (\kappa_\mu/4)|\sin(t^+/200) - \sin(\pi z^+/25)|$, and the overbar denotes an average over a period. The $1/4$ multiple in the latter equation gives the height of the denticles relative to the riblets.

DOLPHIN CHAOS CONTROL

The dolphin’s microgrooves³⁸ are modeled using $\mu^z(t^+) = \mu^x(t^+) = \bar{\mu}(1 + \kappa_{groove} \sin(t^+/50))$, where $\kappa_{groove} = 0.075$. Microvibrations³⁹ are modeled using $I^z_{ext}(t^+) = 0.005(\kappa_{fast} \sin(t^+/800) + \kappa_{slow} \sin(t^+/50))$. Where both microvibration frequencies are applied, $\kappa_{fast} = 1$, $\kappa_{slow} = 1/16$, such that their amplitudes are proportional to the ratio of the periods.

APPENDIX B EXTENDED DATA

This appendix provides additional results that support those given in the main body. It covers:

1. Unsteady boundary layer properties (figure B-1)
2. Streaks (figures B-2 and B-3)
3. Riblet chaos control (figures B-4 and B-5)
4. Shark chaos control (figure B-6)
5. Harmonic riblet chaos control (figure B-7)
6. LEBU (large eddy break-up) device (figure B-8)
7. Stokes' lateral shaking (figure B-9)
8. Temporal phase reset (figure B-10)

UNSTEADY BOUNDARY LAYER PROPERTIES

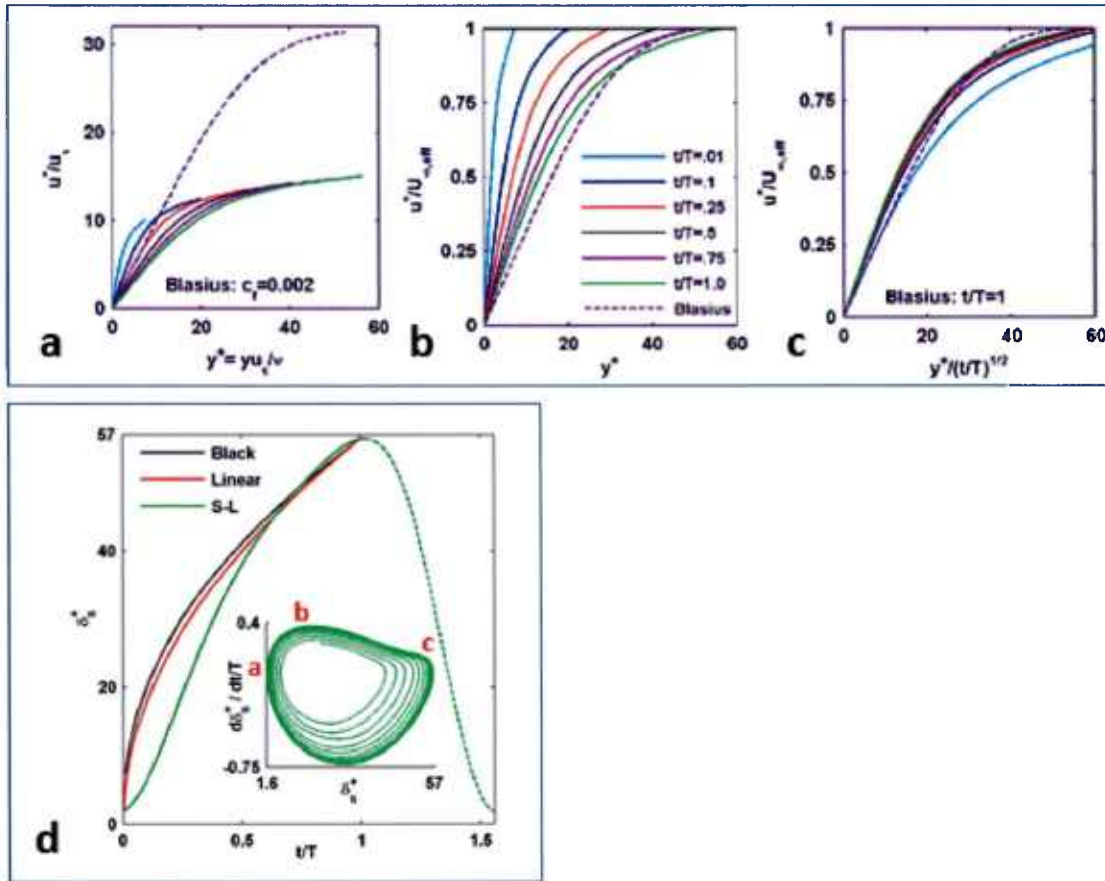


Figure B-1. Viscous Unsteady Sublayer Growth Decay Is a Limit Cycle

a-d: One-dimensional oscillator model of growth and breakdown cycle of sublayer. **a-c:** Unsteady pre-breakdown U -velocity profiles tend to an asymptotic state with $t/T \rightarrow 1$, which is close to the Blasius profile (to get a closed-form solution, the surface-normal component is ignored in the solution of the unsteady momentum equation⁵²). A situation conducive to the formation of TS (Tollmien-Schlichting) waves is approached in each growth-decay cycle (see two-dimensional oscillations in S -triads in figure 3c). **a:** Unsteady velocity profiles (y^+ , u^*/u_τ) obtained by solving NS equations as per Black;⁵² in wall layer coordinates, there is no approach to the Blasius solution. **b:** y^+ , $u^*/U_{\infty,eff}$. (See SI for $U_{\infty,eff}$) **c:** $y^+/(t/T)^{1/2}$, $u^*/U_{\infty,eff}$. **d:** Unsteady sublayer thickness growth ($0 < t/T < 1$) and liquidation ($1 < t/T < 1.5$); linear and SL models are compared with Black's⁵² (t/T , δ_s^+). The linear and SL models do not require any input from the log layer (such as $U_{\infty,eff}$) to reproduce Black's distribution. **d (Inset):** $d\delta_s^+/dt^+$, δ_s^+ . One-dimensional SL model of the sublayer thickness; this model is described in SI. Sublayer thickness variation is a limit cycle varying between y^+ of 1.6 and 57. The sublayer grows during a-b-c, and liquidates the vorticity during c-a. The inset in **d** also shows how a disturbed sublayer would return to the limit cycle once the disturbance is removed; the return path is long (figure B-8b).

STREAKS

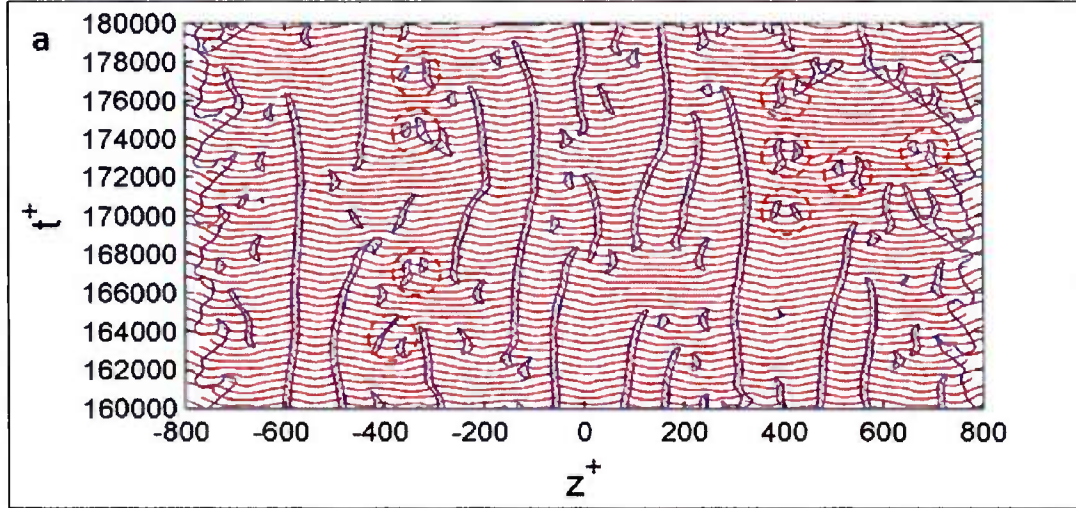


Figure B-2a. Combined Image of Streak and Phase Iso-Contours

The streaks are iso-contours of $|\omega_z^+| = \text{threshold of } A_0/2$. The arrays of wave endings in the phase overlay on the streaks. Corner hairpins (HP): also marked by $|\omega_z^+| = \text{threshold of } A_0/2$, they are leaning forward; there are orthogonal oblique waves travelling along the hairpin length, root to tip, just as they do with streaks. Thus, the longitudinal near-wall vortex pairs and hairpins have similar dynamical properties—orthogonal disturbances travel through them (see wave-guide model of wall-pressure fluctuation in Discussion and Summary). The “noise” between streaks is from local instabilities (waves end or have inflections switching from + to – values); this “noise” does not become streaks and may appear in pairs frequently; therefore, it is not noise. Broken circles: HP lift-up at base of S-wave triads; see figure 4c.¹⁵

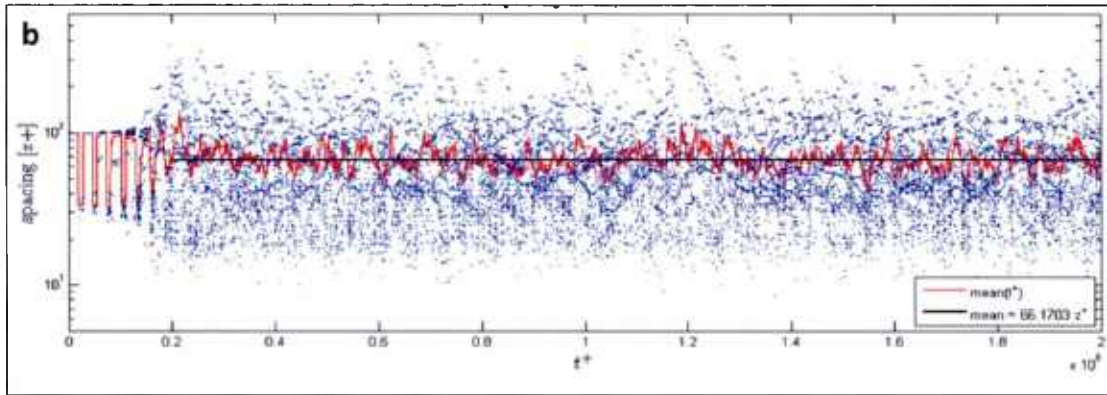


Figure B-2b. Spatio-Temporal Lateral Streak Spacing λ_z^+ in Units of z^+ in Baseline TBL ($t^+ > 0.2 \times 10^4$); Axes (z^+ , t^+)

Lines: Red is moving average at a given t^+ ; black is the mean for $t^+ > 0.2 \times 10^4$. Mean $\lambda_z^+ = 66.17 z^+$, marked by the red circle (figure 2f); this mean compares favorably with experimental data²⁷ (figure 2f); the mean value of λ_z^+ increases if an additional criterion (such as streak length) is imposed, as done in experiments.

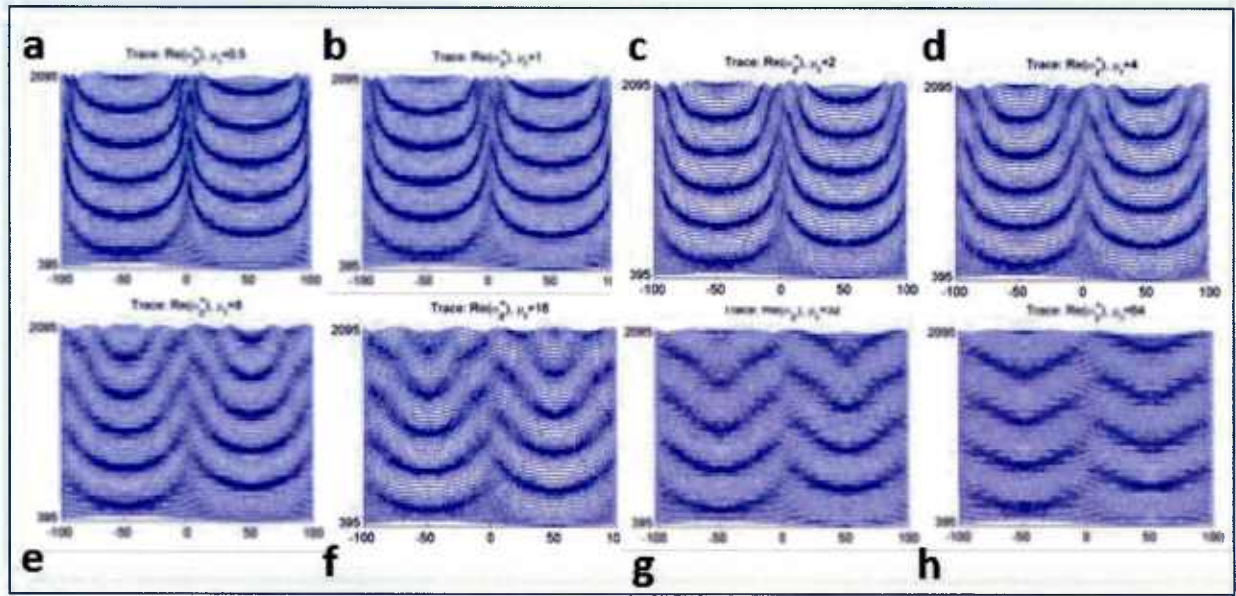


Figure B-3a-h. Effects of Amplitude of μ_r on Lateral Diffusion of Vorticity Waves in TBL

$\mu_r = 0.5$ (a), 1.0 (b), 2 (c), 4 (d), 8 (e), 16 (f), 32 (g), and 64 (h). Number of crashes increases as μ_r drops, and waves approach two-dimensionality as μ_r increases. High Re number TBL has $\mu_r \rightarrow 0$, and unsteady LBL approaching critical Re value has $\mu_r \rightarrow \infty$. Coefficient of lateral diffusion is a measure of Re number, diameter of near-wall vortex pairs and hairpins that form after vorticity is liquidated, and the highness of high-frequency wall-pressure fluctuations (see wave-guide model in figures 2g,h).

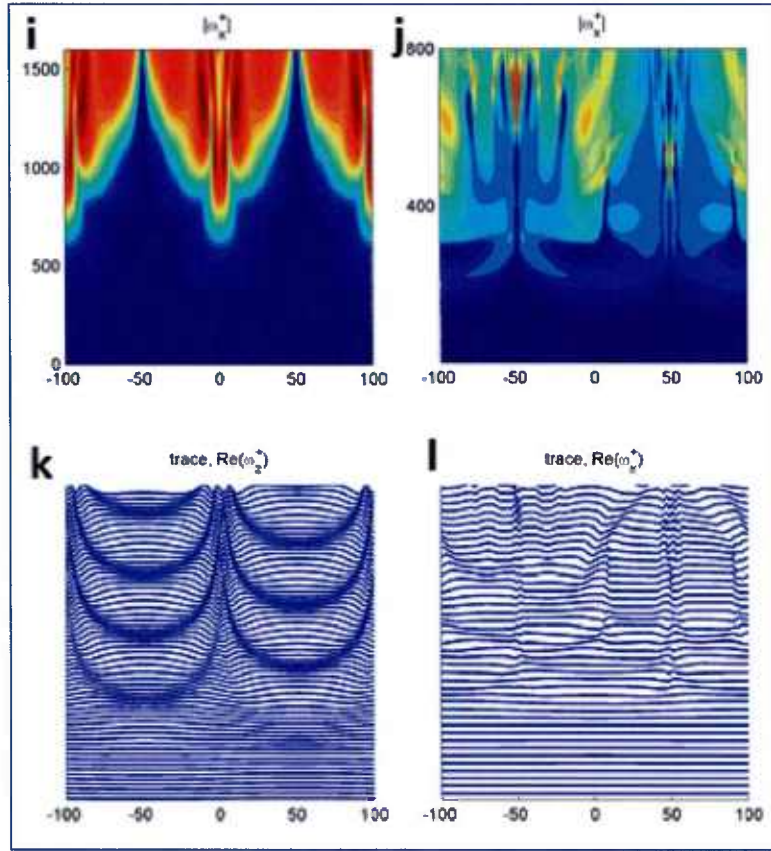


Figure B-3i-l. Laterally Crashing Diffusion Waves in TBL Lead to the Formation of Streaks and Coupling of High and Low Frequencies (see Videos SI1-SI3)

i, k: Coupled ω_z^+ , ω_x^+ oscillators. **j, l:** Coupled ω_z^+ , ω_x^+ , δ_s^+ oscillators. **i, j:** $|\omega_x^+|$ contour. **k, l:** $\text{Re}(\omega_x^+)$. Inclusion of the δ_s^+ oscillator leads to lateral distortion; this three-equation model is described in SI. Lateral crashing is evident in **k** at $z^+ = 0$, where streaks form; this happens at $z^+ = \pm 50$ in **j, l**.

RIBLET CHAOS CONTROL

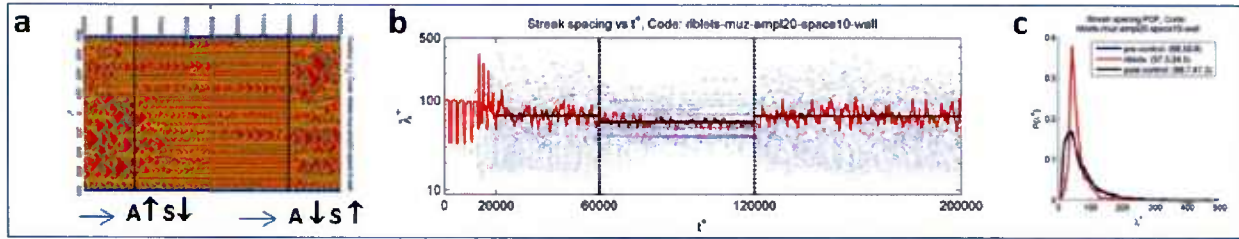


Figure B-4. Shark “Strong” Riblet Model (Excluding Dermal Denticles): Acting on ω_z^+ in the Region Marked by Vertical Dotted Lines in (a, b)

- a:** $|\omega_z^+|$ (z^+ , t^+). “Strong” riblets are active: $60,000 < t^+ < 120,000$, one-equation model (ω_z^+), with “corner” boundary conditions (hairpin shedding in blue); riblet spacing = $10 z^+$.
- b:** Streak spacing in units of z^+ in (t^+ , z^+). Red: Moving average with window. Blue: Individual streak spacing realizations. Black: “Conditional” mean based on whether riblets are active or not. Streak spacing in baseline TBL appears in figure B-2b and figure 2f.
- c:** Probability density function of streak spacing $P(z^+)$, z^+ . Streak spacing is approaching a monochromatic distribution in riblets. In **b**, also notice the parallel blue lines when riblets are applied, indicating that streak spacing is organized into harmonics.

RIBLET CHAOS CONTROL (Cont'd)

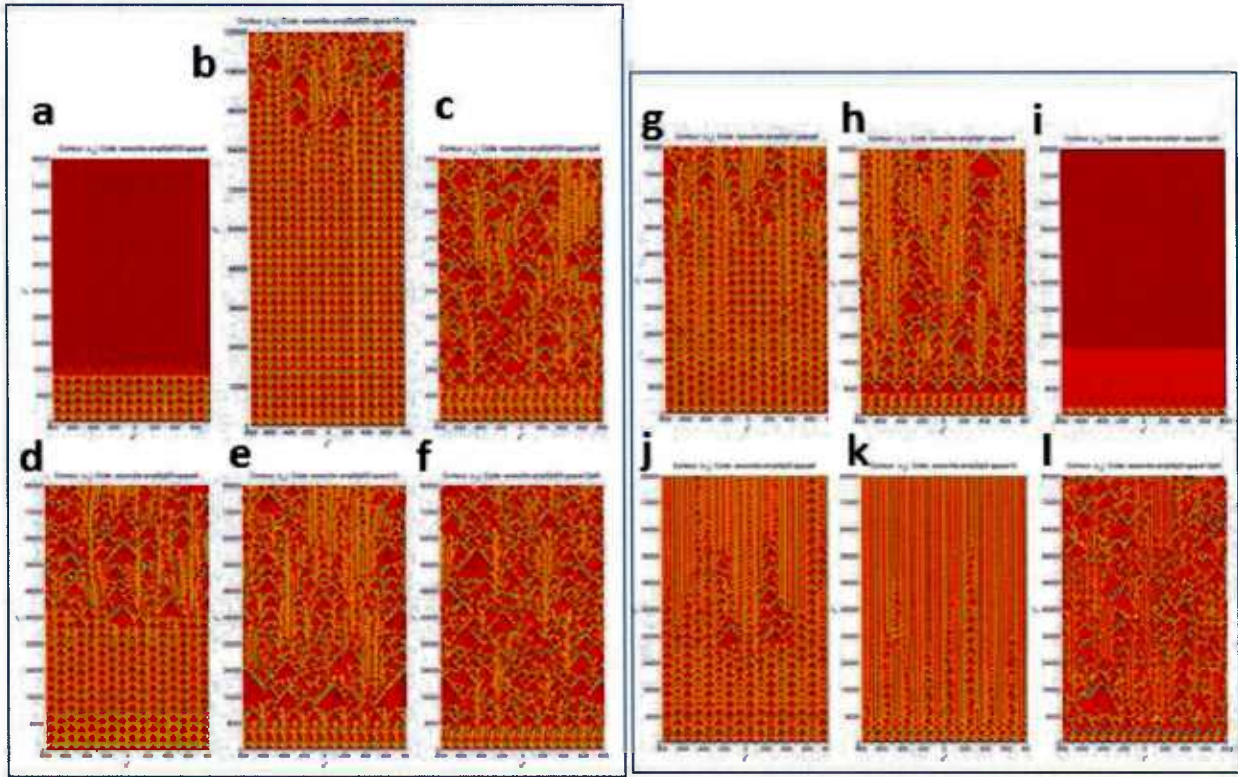


Figure B-5. Riblet Chaos Control

$|\omega_z^+| (z^+, t^+)$; Two-equation (ω_z^+, ω_x^+) model. Effects of diffusion ($\kappa_\mu = 0.025$ (a-c), 0.05 (d-f), 0.1 (g-i), and 0.2 (j-l)) and riblet spacing: $s^+ = 8$ (a, d, g, j), 10 (b, e, h, k), and 12.5 (c, f, i, l). κ_μ, s^+ adjust amplitude and spacing of riblets (see Methods). Baseline TBL: Figure 5f.
b, k: Riblet spacing of $s^+ = 10$ and $\kappa_\mu = 0.025$ (low) and 0.2 (high) give best chaos control, the 0.025 value being for extending transitional range, and 0.2 being suitable when chaos cannot be delayed anymore and TBL has been established.

SHARK CHAOS CONTROL

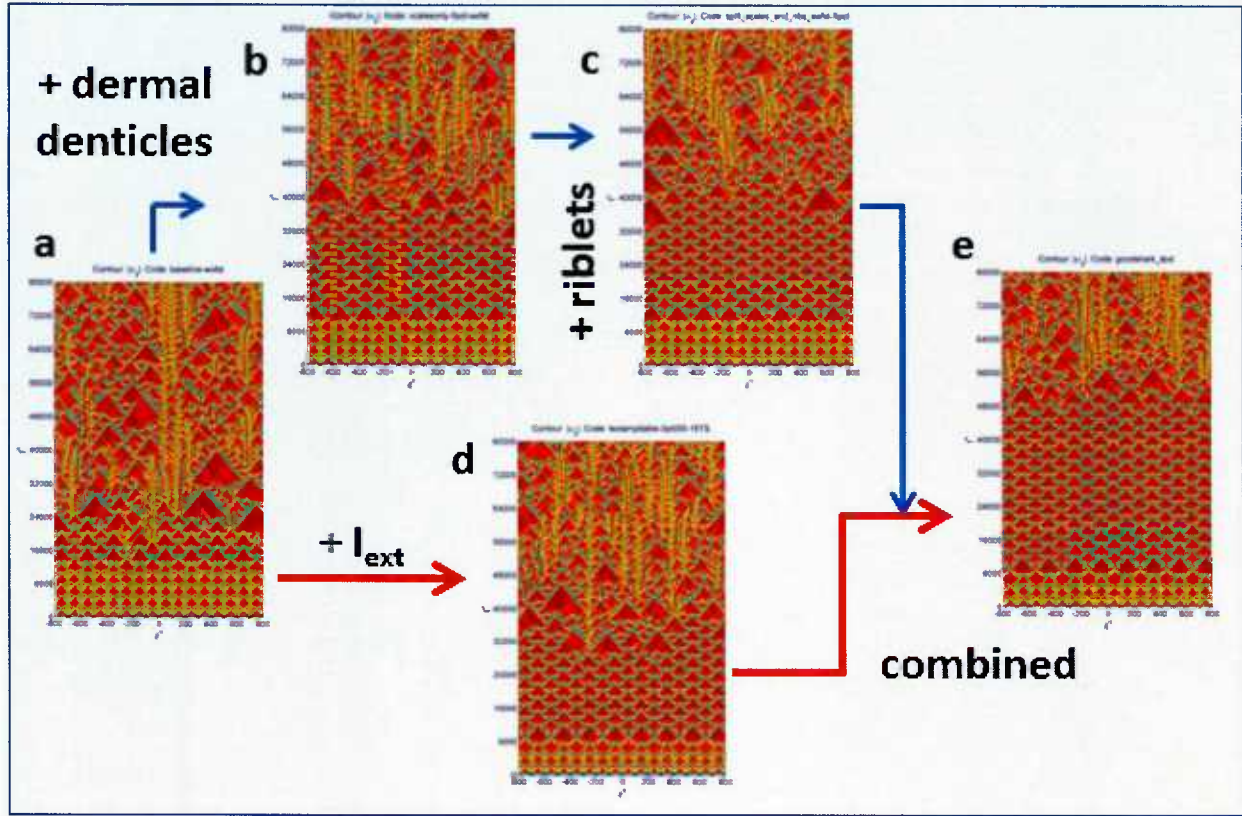


Figure B-6. Chaos Control in Sharks Using Dermal Denticles, Riblets, and Their Vibration

$|\omega_z^+| (z^+, t^+)$. Free-stream turbulence in IC promotes the early onset of chaos. **a:** Baseline TBL. **b:** Effects of dd scales on baseline, where dd scales are imposed on the ω_z^+ oscillator. Onset of chaos is delayed. **c:** Effects of dd and riblets on baseline where riblets are imposed on the ω_x^+ oscillator. Onset of chaos is delayed some more. **d:** External vibration of baseline at a time period of $16^S TS$. **e:** Results of combined dd, riblets, and external stimulation on baseline. Staggered transitional organization is restored, and the onset of chaos is delayed the longest. The lateral oscillators are synchronized. Note that the imposition of the riblets and dd are not precisely the same as used in figure 5; here, $\bar{\mu}^z(z^+, t^+) = \bar{\mu} (1 + Sc(z^+, t^+) - \bar{Sc})$, where $Sc(z^+, t^+) = (\kappa_\mu/4)|\sin(t^+/200) - \sin(\pi z^+/20)|$, and where the overbar denotes an average over a period. $\bar{\mu}^x(z^+, t^+) = \bar{\mu} (1 + \kappa_\mu \sin(2\pi z^+/s^-))$. That is, dd are imposed on the ω_z^+ oscillator, and riblets are imposed on the ω_x^+ oscillator.

HARMONIC RIBLET CHAOS CONTROL

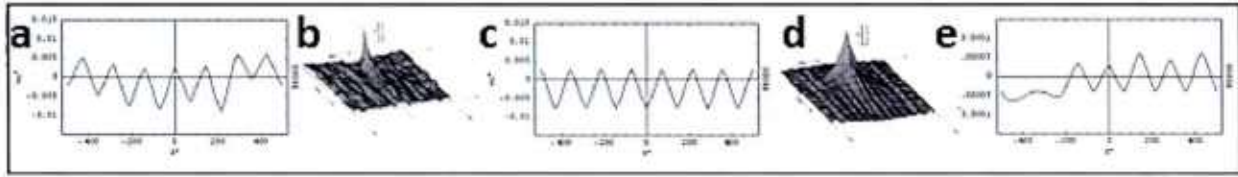


Figure B-7a-e. Alternative Riblet Modeling Using Coupled Oscillators

Present theory is **a**, **c**, and **e**. Riblets are assumed to act as high-frequency oscillators (ω_{xh}^+). **a**, **c**, and **e**: Four-equation model (oscillators ω_z^+ , ω_x^+ , ω_{xh}^+ , δ_s^+) run with low spatial resolution. Details of this model can be found in SI. **a**, **c**: (z^+ , real part of ω_x^+); **e**: (z^+ , real part of ω_{xh}^+); **b**, **d**: Normalized wall-shear stress (t^+ , z^+) due to Choi.⁵³ Animation files of figure B-7a,c,e are provided, with other associated animations, in video SI-5.

Note: Figures B-7b and B-7d are sourced from reference 53.

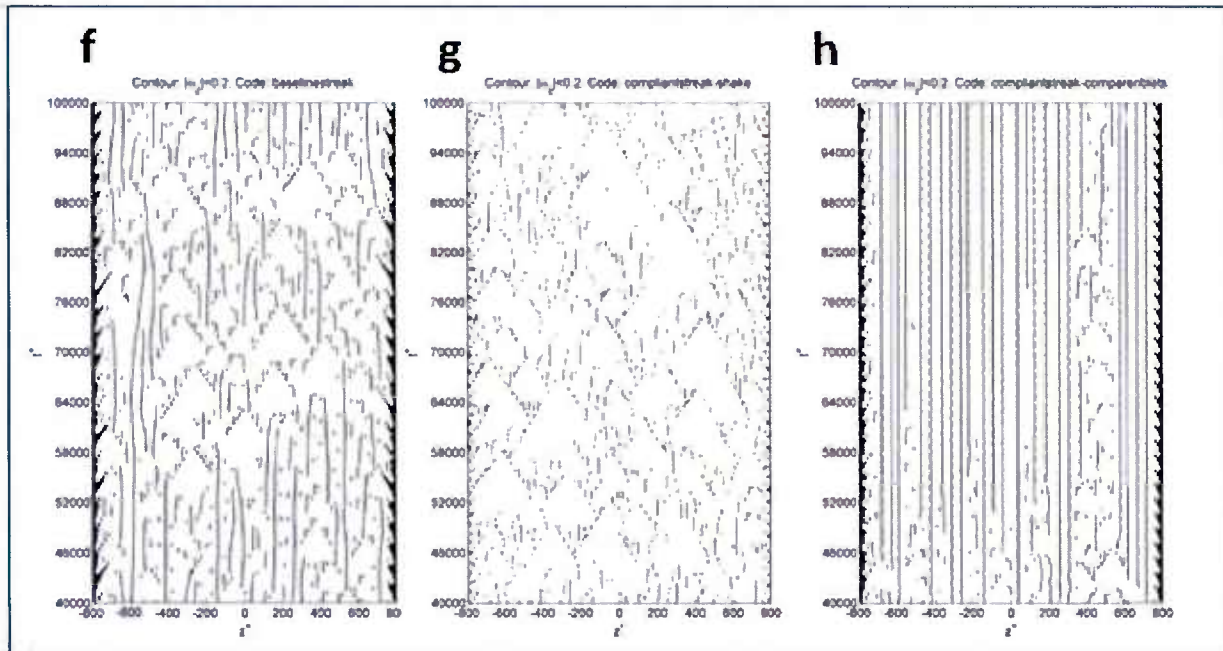


Figure B-7f-h. TBL Streaks Identified by Iso-Contours of $|\omega_z| < 0.2$ in (z^+, t^+)

f: One-equation model (ω_z^+) with corner boundary condition (black hairpins in corners). **g**: Sensing streaks, and applying I_{ext} locally, eliminates streaks. **h**: Strong riblets organize to a preponderance of streaks (compare with figures 5i, 5k, and B-5k, which are from a two-equation model).

LEBU (LARGE EDDY BREAK-UP) DEVICE

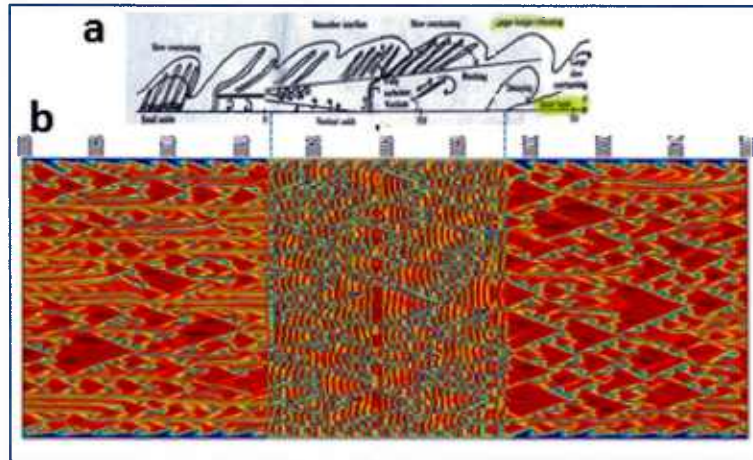


Figure B-8a,b. Chaos Control with Perturbations External to SL Oscillator: LEBU

Present theory is **b**. **a**: TBL hairpin vortices, Kármán shedding from LEBU, outer edge of TBL; x, y ; schematic of smoke flow visualization of TBL perturbed by Kármán shedding from a LEBU;⁴² the region where skin friction is greatly reduced is marked by a pair of vertical broken lines; the near-two-dimensional wake perturbation acts on the TBL between the trailing edge of the LEBU and the onset station of the new internal layer; maximum local skin friction reductions of up to 30 – 40% have been reported; in Stokes' spanwise shaking, up to 50% drag reduction is possible.³⁰ **b**: Present theory: $|\omega_z^+|$ in (z^+, t^+) ; Uncoupled ω_z^+ oscillator ($k_x = 0$), $I_{ext} = A_0 \sin(2\pi t^+/T)$ (method appears in SI); TBL returns with a vengeance. In an animal, this would be the tail region, which would keep the tailcone TBL attached. See video SI-4.

Note: Figure B-8a is sourced from reference 42.

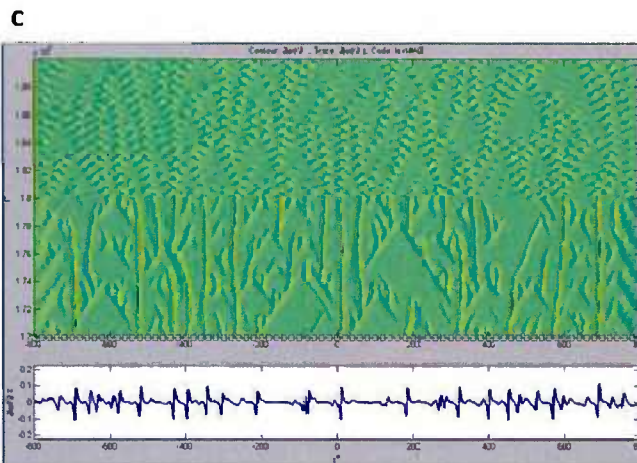


Figure B-8c. External Perturbation I_{ext} Breaks Streaks

Streaks are detected by $\partial|\omega_z|/\partial z$, as shown in lower trace at the cursor location (horizontal chain line in upper trace). See video SI-4.

STOKES' LATERAL SHAKING

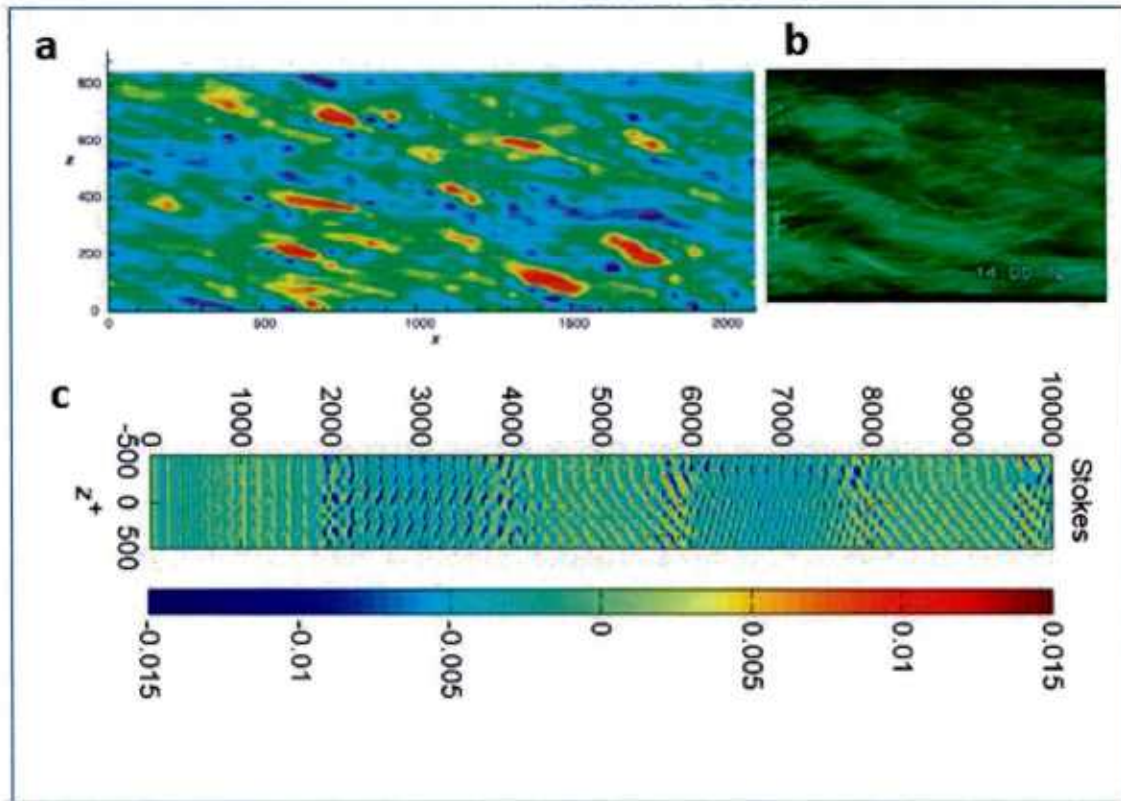


Figure B-9. Chaos Control with Perturbations External to SL Oscillator: Stokes' Drag Reduction

a: Numerical simulation of streamwise velocity contour;⁵⁴ spanwise oscillatory force control on TBL. **b:** Smoke tracer concentration distribution.⁴³ **c:** This theory; the $|\omega_z^+| (z^+, t^+)$ effects of Stokes' lateral wall shaking on drag reduction is modeled as I_{ext} . The three-equation model that yields the results of **c** are provided in SI, with accompanying justification of the terms employed.

Note: Figure B-9a is sourced from reference 54. Figure B-9b is sourced from reference 43.

TEMPORAL PHASE RESET

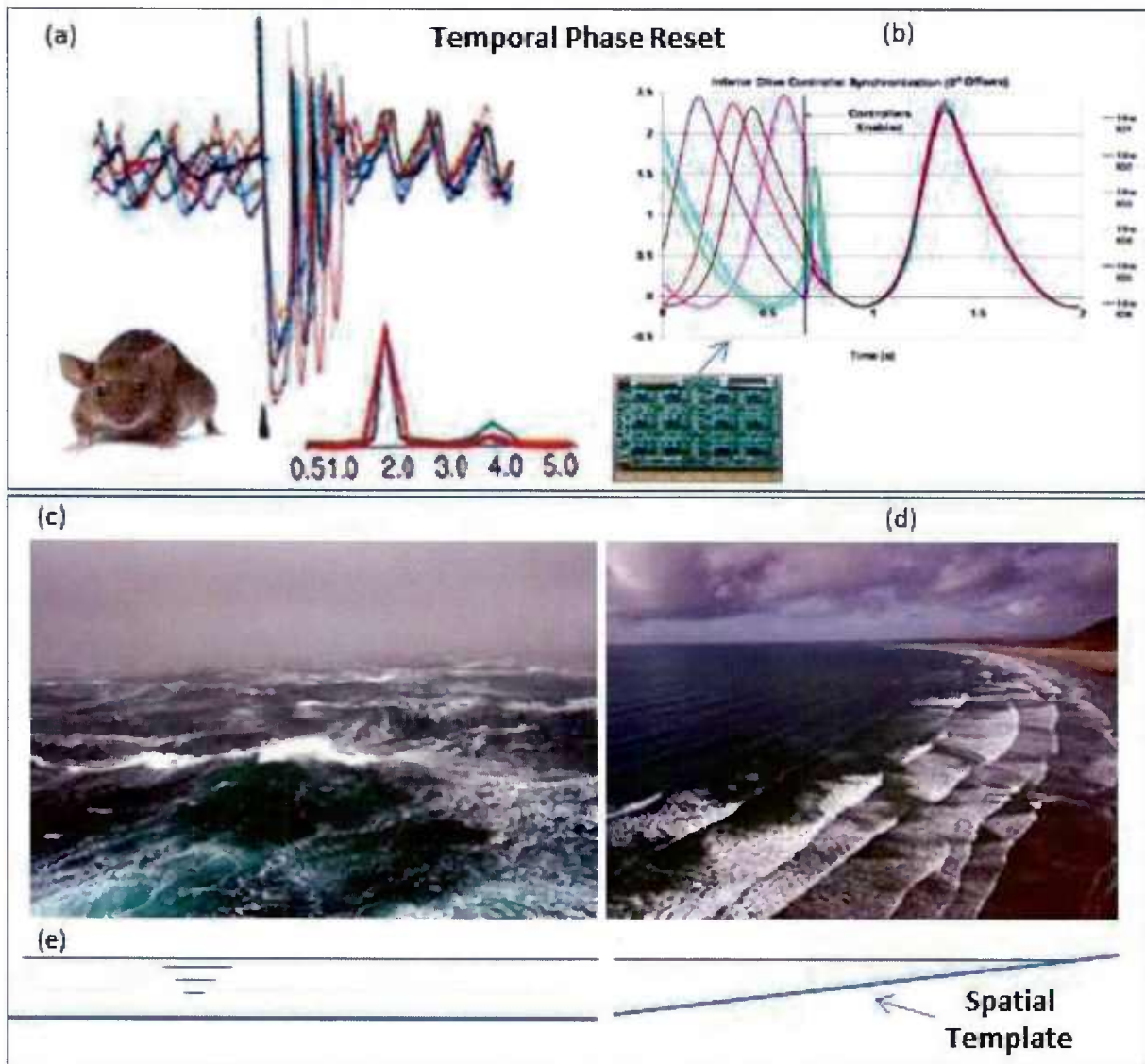


Figure B-10. Schematic Explaining Spatio-Temporal Phase Reset

a: Impulse causes phase reset in 10 rodent brain inferior-olive (IO) neuron time traces.^{1,55}

b: Analog solutions of six IO neurons (lower inset in **b**) are synchronized as in **a**.²³

c-e: A spatial template (**e** right: sloping shore) makes the deep ocean (**c** left) waves⁵⁶ (**c**) become two-dimensional.⁵⁷ **d:** *a, b* are temporal; *c, d* are spatio-temporal.

Note: Reference sources: Figure B-10a is from reference 1; B-10b is from reference 23; rodent in B-10a is from reference 55; photographs in B-10c and B-10d are from references 56 and 57, respectively.

APPENDIX C SUPPLEMENTARY INFORMATION (SI)

SI METHODS AND DISCUSSION

This appendix describes the methods used to generate the figures provided in the extended data of appendix B. It also discusses these methods and provides further discussion of the results in the main body. It covers:

1. Unsteady boundary layer properties:
 - Linear sublayer thickness model
 - Log-layer bounded sublayer thickness model
 - Stuart-Landau sublayer thickness model
 - Inclusion of sublayer thickness in model: Coupling between vorticity and sublayer equations.
2. Shark chaos control
3. LEBU (Large Eddy Break-Up) device
4. Stokes' lateral shaking for drag reduction

VIDEOS

Note: Videos are available from the principal author.

Other supplementary information available for download includes the videos listed below. These videos are animations that make certain results more clear:

Video SI-1. Animation of the ω_z^+ laterally crashing waves using the model shown in still form in figure 4f. It shows how streaks are formed. Oscillators: ω_z^+ as per baseline (wavelength of 200 wall units; Methods), $\omega_x^+(z^+, 0) = 0$, rather than noise.

Video SI-2. Animation of $Re(\omega_z^+)$ crashing waves using ω_z^+ initial condition wavelength $\lambda_{IC} = 200z^+$. Clear formation of streaks and carrier waves when $\lambda_{IC} = 200z^+$, compared to $100z^+$. Streaks form where $Re(\omega_z^+)$ crosses sign with spacing $\lambda_{IC}/2$. Oscillators: ω_z^+ , ω_x^+ , δ_s^+ ; IC: $\omega_x^+(z^+, 0) = 0$, $\delta_s^+(z^+, 0) = 10$; model details given later in this appendix.

Video SI-3. Animation of $Re(\omega_z^+)$ crashing waves using initial condition wavelength $\lambda_{IC} = 100z^+$. Compare crashing behavior to Video SI-2; crashing is not as clear when initial wave

length is 100 wall units. Oscillators: ω_z^+ , ω_x^+ , δ_s^+ ; IC: $\omega_x^+(z^+,0) = 0$, $\delta_s^+(z^+,0) = 10$; model details given later in this appendix.

Video SI-4. Animation of $\partial|\omega_z^+|/\partial z^+$, where contours and traces of that variable are plotted simultaneously with a moving cursor. A strong oscillator (LEBU) is applied at $t^+ = 180,000$. This video shows how streaks are annihilated. Also see the stationary view in figure B-8b.

Video SI-5. This is a presentation-like movie containing animations of ω_z^+ , ω_x^+ , and ω_{xh}^+ , for cases with and without harmonic riblets. Increasing lock-in with increasing number of control oscillators is apparent. See figures B-7a,c,e for stationary view and see this appendix for accompanying methods and discussion.

UNSTEADY BOUNDARY LAYER PROPERTIES

In figure B-1a, the profiles match the log law and $U^+ = y^+$ equation at the two extremities of y^+ ; this entire region is viscous; $U_{\infty,eff}$ is unsteady and is 99% of the velocity obtained by extrapolating the log layer wall-ward at the given t/T .⁵² In figure B-1c, the horizontal axis $y^+/\sqrt{t/T}$ is reminiscent of Blasius boundary layer thickness scaling of $\eta = y\sqrt{\frac{U_\infty}{\nu x}}$, if $x/t = U$, a velocity, and ν = kinematic viscosity.⁵⁸ Instead of the local value, a value of the friction velocity is used that is taken as the average over the cycle in B-1d (an attractor value⁵⁹). In figure B-1d, the vorticity in the layer $y^+ < 1.6$ is not liquidated;²⁹ not liquidating down to the wall satisfies the requirement of dynamical separation;⁶⁰ see p. 302 of Bandyopadhyay et al.⁶¹ and also Bandyopadhyay.⁵⁹

The viscous boundary layer results in figure B-1 can be examined in the context of unsteady laminar boundary layer theory.⁴⁴ Consider the unsteady laminar problem of oil emanating into quiescent fluid from a tube at $r = 0$ and spreading radially (or two-dimensionally) over a surface such that the leading edge is at the point $r = R(t)$, producing a boundary layer in the quiescent fluid. Similarity solutions of the unsteady momentum and continuity equations asymptotically reduce to the Blasius relationship as $r \rightarrow R$. In the context of figure B-1d, the beginning of each cycle (marked **a**) is analogous to the region in which the oil spill's leading edge is in the region of $r = 0$, although this exact point is never reached because vorticity is not liquidated all the way to $y^+ = 0$.^{29,60} The velocity profiles in the present model approach the Blasius solution with increasing t/T , which is analogous to Phillips' region $R(t) \gg 0$; in Phillips' theory, the oil film is stretched as per **a-b-c**, and the Blasius profile is approached at **c**, whence TS waves (named pseudo^sTS here) likely form as per the present model. Einstein & Li's observation²⁶ that the viscous layer growth in a TBL slows as the layer thickness grows (region **b-c** in the figure B-1d inset) indicates that the layer is within the $R(t) \gg 0$ region—in which the Blasius relationship is closely approached—over a large fraction of each self-regulating cycle (the region **a-b** is like $r \rightarrow 0$ and the region **b-c** is like $r \rightarrow R$). The “oil-spill” cycle is repeated as per the self-regulation cycle. The present model has been put in the context of Phillips' more rigorous theory and beyond Black's because it shows that a Blasius solution is indeed approached asymptotically, whereby the idea of^sTS-waves forming within the TBL during each production cycle is reasonable.

Linear Sublayer Thickness Model

In the linear model in figure B-1d, the evolution of the sublayer thickness δ_s^+ is modeled by the expression

$$\dot{\delta}_s^+ = \beta \frac{1}{\delta_s^+} - \text{liquidation.} \quad (\text{C-1})$$

The first term on the right-hand side of the expression is the growth term; the liquidation term is not used in the curves in figure B-1d and will be discussed later. The form of the growth term is informed by the work of Einstein and Li,²⁶ who solved the diffusion equation

$$\frac{\partial u}{\partial t} = \nu \frac{\partial^2 u}{\partial y^2}$$

with the initial and boundary conditions

$$u(0, y) = U_{ref} \quad u(t, 0) = 0 \quad \lim_{y \rightarrow \infty} u(t, y) = U_{ref},$$

which yields the expression for the time-dependent velocity profile within the sublayer:

$$u(t, y) = U_{ref} \operatorname{erf}\left(\frac{y}{2\sqrt{\nu t}}\right).$$

In the above expressions, the scaling velocity U_{ref} is a velocity that exists at some representative height within the turbulent flow. Einstein and Li's final expression can be manipulated to find a time-dependent expression for the thickness of the sublayer δ_s , defined as 99% of U_{ref} :

$$\delta_s = 2\sqrt{\nu t} \operatorname{erf}^{-1}(0.99).$$

The above expression may be recast in wall units⁵⁸ such that

$$\delta_s^+ = 2 \operatorname{erf}^{-1}(0.99) \sqrt{t^+}.$$

The time derivative of the above expression yields

$$\dot{\delta}_s^+ = \frac{\operatorname{erf}^{-1}(0.99)}{\sqrt{t^+}} = \frac{2 (\operatorname{erf}^{-1}(0.99))^2}{\delta_s^+} \approx \frac{6.635}{\delta_s^+},$$

i.e., the form of the growth term in equation (C-1), such that $\beta = 6.635$.

Log-Layer Bounded Sublayer Thickness Model

Black⁵² also uses the diffusion equation as the basis for his sublayer thickness model but imposes initial and boundary conditions:

$$\begin{aligned} u(y, 0) &= u_\tau \left(A + B \ln \left(y \frac{u_\tau}{\nu} \right) \right), & \text{for } y > 0 \\ u(0, t) &= 0, & \text{for } 0 < t \leq T \\ \lim_{y \rightarrow \infty} u(y, t) &= u_\tau \left(A + B \ln \left(y \frac{u_\tau}{\nu} \right) \right), & \text{for } t > 0 \end{aligned}$$

where A and B are experimentally determined parameters for the log layer, and u_τ is the friction velocity. The expression (see equation 2.31 in Black⁵²) for the velocity $u^* = u/u_\tau$ is a function of the non-dimensional time t/T . The sublayer thickness is the height at which u^* is within 1% of the log-layer's value at that time; the effective free-stream velocity $U_{\infty,eff}$ is the velocity at the sublayer's outer boundary.

Stuart-Landau Sublayer Thickness Model

A modified SL expression can instead be used to model the thickness of the sublayer:

$$\delta_s^+ = \left(\frac{\beta}{\Re(\delta_s^+)} + \sigma \delta_s^+ - \frac{\lambda}{d_\delta} |\delta_s^+ - \kappa|^2 (\delta_s^+ - \kappa) \right),$$

where β and the SL constants σ and λ are as described earlier in this section and in Methods, respectively. The coefficients d_δ and κ are chosen together to yield a limit cycle approximating that of the linear and Black models; in figure B-1d, $d_\delta = 2400$ and $\kappa = 53 + 20i$. Note that in the absence of a liquidation term, the linear and log-layer-bounded models indicate a sublayer growing without bound, whereas the SL model is intrinsically self-regulating.

Inclusion of Sublayer Thickness in Model: Coupling Between Vorticity and Sublayer Equations

Figures B-3(j,l), B-7(a,c,e), and B-9c use vorticity equations that are coupled to the linear model of the sublayer thickness, such that

$$\dot{\omega}_z^+ = \left(\sigma \omega_z^+ + \mu \frac{\partial^2 \omega_z^+}{\partial z^{+2}} - \lambda^z \omega_z^{+2} \omega_z^{+*} \right) k_\delta \delta_s^+ - \hat{k}_x |\omega_x^+|, \quad (\text{C-2a})$$

$$\dot{\omega}_x^+ = \left(\sigma \omega_x^+ + \mu \frac{\partial^2 \omega_x^+}{\partial z^{+2}} - \lambda^x \omega_x^{+2} \omega_x^{+*} \right) k_\delta \delta_s^+ - \hat{k}_h |\omega_{xh}^+| + \hat{k}_z \frac{\partial \omega_z^+}{\partial z^+}, \quad (\text{C-2b})$$

$$\dot{\omega}_{xh}^+ = \left(\sigma \omega_{xh}^+ + \mu \frac{\partial^2 \omega_{xh}^+}{\partial z^{+2}} - \lambda^h \omega_{xh}^{+2} \omega_{xh}^{+*} \right) \tau_h k_\delta \delta_s^+ + \hat{k}_{xh} \frac{\partial \omega_x^+}{\partial z^+} \quad (\text{C-2c})$$

$$\delta_s^+ = \beta \left(g_0 + \frac{1}{\delta_s^+} \right) - k_l \frac{|\Re(\omega_z^+)| [\text{sgn}(\Re(\omega_z^+) - \Delta) + 1]}{2}, \quad (\text{C-2d})$$

where the sublayer liquidation term, omitted in equation (C-1), has now been included. In the expressions for the spanwise and streamwise vorticity, the time scale $k_\delta \delta_s^+$ has physical meaning in that vorticity perturbations may be assumed to grow more quickly in a thicker sublayer. The coupling between ω_z^+ and ω_x^+ is assumed for now to be “outside” this sublayer-based time coupling. The coefficients \hat{k}_x and \hat{k}_z are modified from the values provided in Methods (appendix A) based on order-of-magnitude considerations such that $\hat{k}_x = 0.1 k_\delta \bar{\delta}_s^+$, where $\bar{\delta}_s^+ \approx 30$ is the mean thickness of the sublayer over a cycle, and the multiplier $k_\delta = 0.00025$. Unless noted, the value $\mu_r = 0.49$ is used when the above model is used, instead of $\mu_r = 16$ as described in Methods. This lower value of μ_r is indicative of a higher Reynolds number (figure B-3).

An additional equation for the harmonic vorticity ω_{xh}^+ , which was used in an early iteration of riblet modeling (figures B-7a,c,e), has also been included. When no harmonic riblets are applied, setting $\hat{k}_h = 0$ is sufficient to revert this four-equation model to a three-equation model in $\omega_z^+, \omega_x^+, \delta_s^+$; the terms “three-equation model” and “four-equation model” are used to refer to sublayer-coupled models run, respectively, without and with harmonic riblets. *The physical justification of this riblet model, which here is called harmonic riblets to distinguish it from the aspect ratio-based modeling described in Methods, is that individual riblets were proposed to create local vortices of their own scale, which were driven by the (small) spanwise velocity near the wall.* These local vortices can be seen in visualizations.⁶² The reader will note that this oscillator is coupled to the ω_x^+ oscillator in the same way that the ω_x^+ oscillator is coupled to ω_z^+ , and the respective constants \hat{k}_h and \hat{k}_{xh} match \hat{k}_x and \hat{k}_z , when the harmonic oscillator is used. A relative time constant τ_h has been added to this equation as well, such that these near-wall control oscillators can be given different frequencies from the “primary” three oscillators. $\tau_h = 8$ has been used in the present work.

The above form of coupling and liquidation in equation (C-2d) performs well at low spatial resolution, where the resolution imposes a kind of averaging on the calculations. However, at higher spatial resolution, this sublayer model has a tendency to yield unphysical “spires” of unliquidated sublayer as simulation time increases. The sublayer coupled results in figures B-3j and 3l (right column) are shown at low simulation time. In figures B-7a,c,e, the spatial resolution is low ($dz^+ = 142$), and the simulation could be continued for high values of t^+ .

The constant growth term g_0 added to the expression for sublayer thickness has been estimated as the proportion of the growth layer that Popovich and Hummel²⁹ found was not liquidated, such that $g_0 \approx 0.032$. The liquidation threshold is $\Delta = k_\Delta A_0$, where sublayer intermittency $\gamma \approx 1 - k_\Delta$. Here, sublayer intermittency is defined as the proportion of time for which the sublayer is being actively liquidated. It has been taken as $\gamma = 0.1$, such that the liquidation time is very short, relative to the growth time.⁵²

The liquidation coefficient k_l can be estimated from the values of other constants since the net growth of the sublayer over a single growth-breakdown cycle is small. Equation (C-2d) can therefore be integrated over a cycle of length T to give the condition

$$\int_0^T \beta \left(g_0 + \frac{1}{\delta_s^+} \right) dt^+ - \int_0^T k_l |\Re(\omega_z^+)| \frac{(\text{sgn}(\Re(\omega_z^+) - \Delta) + 1)}{2} dt^+ \approx 0.$$

Since the second integral is zero when the condition $\Re(\omega_z^+) < \Delta$ holds, the above equation may be rewritten in the form

$$\int_0^T \beta \left(g_0 + \frac{1}{\delta_s^+} \right) dt^+ \approx \int_0^{T^\gamma} k_l |\Re(\omega_z^+)| dt^+. \quad (\text{C-3})$$

The following approximations are introduced:

$$\delta_s^+ \approx \sqrt{2\beta t^+}, \quad \Re(\omega_z^+) \approx k_\Delta A_0, \quad \gamma \approx (1 - k_\Delta).$$

Equation (C-3) may now be evaluated to find the approximation

$$k_l \approx \frac{\beta}{A_0(1-k_\Delta)k_\Delta} \left(g_0 + \left(\frac{2}{\beta T} \right)^{\frac{1}{2}} \right).$$

Unfortunately, the time period T is not known *a priori*; instead, the approximation

$$T = \frac{2\pi}{\omega_{base} k_\delta \delta_{s,ref}^+}$$

has been used, where the base frequency $\omega_{base} = c_0 - c_2$. The value of the reference sublayer thickness $\delta_{s,ref}^+$ is taken to be the minimum sublayer thickness $\delta_{s,min}^+ = 1.6$, taken from Popovich and Hummel.²⁹ To account for the imprecision in the approximations, the expression

$$k_l = \hat{k}_l \frac{\beta}{A_0(1-k_\Delta)k_\Delta} \left(g_0 + \sqrt{\frac{\omega_{base} k_\delta \delta_{s,min}^+}{\pi \beta}} \right)$$

has been employed, where \hat{k}_l is a coefficient near one. Note that the method used to approximate the period T is an overestimation since $\delta_s^+ \geq \delta_{s,min}^+$, meaning that in general $\hat{k}_l \geq 1$.

SHARK CHAOS CONTROL

Figures B-4b and B-4c show that, due to riblet application, streak spacing approaches a monochromatic spacing.

Figure B-5 shows that between riblet spanwise spacing of 8, 10, and 12.5 wall units, the optimal spacing is 10 wall units. Chaos onset is best delayed at κ_μ of 0.025. When chaos onset cannot be prevented, then an amplitude of 0.20 is required for best organization.

In figure B-5, the best chaos control (arrayed K-type waves) is achieved at $\kappa_\mu = 0.025$ (weak diffusion), and riblet spacing of 10 wall units. With increasing κ_μ , at $s^+ = 10$, the arrayed K-type

transitions to arrayed A-type in k. The t^+ limit is 120×10^3 in figure B-5b, whereas it is 80×10^3 in the rest. The S- and A-waves also display a subdued tendency, oscillating between streaks with restrained tendency for streak jumps and growth in S-waves. *The riblets have brought in-phase balance between the activator and inhibitor parts²¹ of the self-regulating instability (the positive and negative feedbacks are in phase).*

Normally in a TBL, the wave triads are chaotically disposed. However, figures 5 and B-6 show that the shark dd's and their staggering provide a template for the TBL instabilities to lock in to because the templates mimic the primary modes of the flow. The lock-in mechanism is a spatio-temporal version of the temporal Self-referential Phase Reset (SPR).^{1,23} Although the process is nonlinear, chaos is eliminated and the initial condition dependence is removed by the effects of the template, which acts as an impulse of required amplitude and duration.

The dependence on IC can also be removed (figures B-7a,c,e) by imposing the harmonic riblets model described earlier. In such a model, $\frac{\lambda^z}{\lambda^x} = \frac{\lambda^x}{\lambda^h} = \frac{1}{100^2}$ (figures B-7a,c,e). The baseline TBL spanwise ω_z^+ vorticity distributions have both standing and transverse oscillations (figure B-7a). When a weak orthogonal streamwise vorticity ω_x^+ oscillator (ω_{xh}^+) is added, the transverse oscillations in ω_z^+ are reduced and they mostly stand and oscillate (pure varicose instability). However, when the riblet oscillator (ω_{xh}^+) is added, the ω_z^+ pattern standing becomes locked and all transverse waves are eliminated (figure B-7c). The lock-in is qualitatively in agreement with measurements.⁵³ The distributions of ω_{xh}^+ (figure B-7e) show that the riblets apply a correction to the other oscillators to achieve aligned lock-in, but this correction is applied by only some of the riblets at a time; that is, ***at a given instant of time, only some of the riblets are active*** (figure B-7e). Note that the different methods of riblet modeling (surface template or harmonic) share this quality, which is also observed experimentally (figure 10 of Lee and Lee⁶²). For temporal representation of figure B-7, see video SI-5.

Paraphrasing the stability theory due to Phillips,⁶³ the streamwise ω_x^+ oscillator couples with the ω_z^+ transverse and main oscillators, which leads to the formation of sub-streak wavelengths (Phillips' figure 4); Phillips' theory hints that the riblet amenable sub-streak wavelength is naturally present but submerged in the chaos of streak spacing in figure B-2a; once the riblet template (which in figure B-7c is ω_{xh}^+) is offered, this wavelength locks-in to that—a Llinas¹ phase reset takes place.

The wing-tip to wing-tip spacing of dermal denticles of sharks is on the order of streak spacing (figures 2f and B-2b). The dd surface pattern would help organize the streak spacing (figure B-6b), which in turn would organize the sub-streak wavelengths of Phillips, discussed above. As a result, the shark riblets would have less work to do.

In figure B-7a, both surface-normal and lateral travelling waves are present and the distribution is similar to that measured by Choi⁵³ (figure B-7b is a reproduction of figure 17 from Choi), where spanwise correlation of skin friction is roughly *stationary and non-uniform*. However, in figure B-7c, the pattern is purely *stationary and uniform* and can be compared with measurements in figure B-7d.⁵³

Figures B-7a,c,e give theoretical confirmation to Watson et al.,⁶⁴ who noted that (italics added by present authors):

Turbulent boundary layers exhibit spanwise variations in skin friction coefficients and integral boundary layer properties, even in flat plate experiments where great care has been taken to insure the two-dimensionality of the outer flow. Spanwise plots of [momentum thickness] for nominally two-dimensional flows are always somewhat ragged, as are [skin friction] measurements. *The reason may be because of the raggedness in the transition process, and/or the presence of near wall structures in the boundary layer, inducing a standing pattern in the flow.*

Considering the induced wall image, the standing pattern is interpreted as varicose instability; they are rings in B-7c and spirals in B-7a (see figure 1 in Desyatnikov et al.,⁶⁵ who solve the nonlinear Schrödinger equation and interpret similar results as solitons of self-trapped waves; this is in broad agreement with the present conclusion that streaks are solitary waves trapped in the near-wall viscous layer).

The vorticity pattern in figure B-7c oscillates in the vertical direction while remaining stationary in the transverse direction. This shows the presence of varicose instability (video SI-5). (In the presence of a transverse wave, the vorticity layer spirals left or right.) From this, the liquidation of vorticity into hairpin vortices may be considered to be analogous to volcanic eruption due to plate tectonics. (The plate becomes locally thin and is unable to prevent the high-pressure hot magma from erupting.) The vorticity layer at a given (z^+, t^+) thickens and thins due to stretching in two orthogonal directions parallel to the wall. If a triad is situated there and is resonating as in S-type, for a brief instant the layer thickness suddenly reaches the critical height in all three TS waves simultaneously (in arrayed A-triads, it doesn't). At that instant, the diffused vorticity spins faster while thinning, becomes concentrated (discrete), develops amplifying local oscillation briefly, bifurcates (selects handedness²⁴), and subsequently pinches off into the recognizable hairpin form. Due to Biot-Savart wall imaging, the vortex line lifts up into a hairpin form.

Figure B-7g, in comparison to B-7f, shows that when the long streaks ($|\omega_z^+| < 0.2$ contours) in the baseline TBL are detected and shaken (using some I_{ext}), they are eliminated. However, when the short oblique “streaks” are shaken, they become fragmented.

Figure B-7h shows that when a strong lateral diffusion perturbation is added to the TBL to model riblets at spacing $z^+ = 10$, the TBL becomes populated with rows of parallel streaks; this A-type dominant pattern is also seen in figures 5i,k.

LEBU (LARGE EDDY BREAK-UP) DEVICE

The model in figure B-8 is as follows. The powerful external oscillator used to model the effect of vortex-shedding devices placed into the outer layer is given by the expression

$$I_{ext}^z(t^+) = A_0 \sin\left(\frac{2\pi t^+}{1000}\right).$$

Note that this period of oscillation is not an even multiple of T_{TS} . This oscillator's amplitude is equal to that of the most powerful wall vorticity oscillator ω_z^+ , and so may be termed a “very strong” perturbation, as might be expected if coherent vortices produced with higher speed outer-layer fluid are brought close to the wall. The rest of the model is a one-equation model for ω_z^+ (i.e., $k_x = 0$), as given in Methods.

STOKES' LATERAL SHAKING FOR DRAG REDUCTION

Stokes' method of drag reduction (shown in figure B-9), imposed experimentally by oscillating the wall in the spanwise direction, is modeled using an external perturbation to equation (C-2b) of the sublayer-coupled model, such that

$$\dot{\omega}_x^+ = \left(\sigma \omega_x^+ + \mu \frac{\partial^2 \omega_x^+}{\partial z^{+2}} - \lambda^x \omega_x^{+2} \omega_x^{+*} \right) k_\delta \delta_s^+ + \hat{k}_z \frac{\partial \omega_z^+}{\partial z^+} + I_{ext}^x(t^+),$$

$$I_{ext}^x = A_0 \frac{\lambda_r^z}{\lambda_r^x} \times \frac{1}{10} \sin(\omega_{base} k_\delta \delta_{s,min}^+ t^+) = A_0 \frac{1}{1000} \sin\left(\frac{t^+}{625}\right).$$

Physically, this perturbation corresponds to an external source of streamwise vorticity, $\omega_x = \partial w / \partial y - \partial v / \partial z$, where v, w are the surface-normal and spanwise velocities, respectively. When imposed in a highly sheared TBL, Stokes' layer does not penetrate as far from the wall as it does in an unsheared fluid,³⁰ meaning that at some (small) height, the spanwise velocity is virtually unperturbed by the oscillations of the wall. Thus, for $v \approx 0$ (non-penetrating condition), it is permissible to say that a velocity condition $w(0, t) = f(t)$ is equivalent to a source of streamwise vorticity.

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